Brownian motion (cont.)

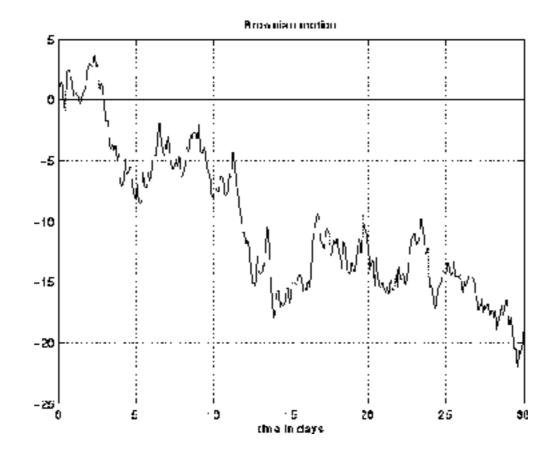
18.S995 - L05

1.2 Brownian motion

Diffusion equation with constant drift

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Path-wise representation of typical trajectories?



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1.2.1 SDEs and discretization rules

The continuous stochastic process X(t) described by Eq. (1.19a) or, equivalently, Eq. (1.20) can also be represented by the stochastic differential equation

$$dX(t) = u dt + \sqrt{2D} dB(t). \tag{1.25}$$



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Here, dX(t) = X(t + dt) - X(t) is increment of the stochastic particle trajectory X(t), whilst dB(t) = B(t + dt) - B(t) denotes an increment of the standard Brownian motion (or Wiener) process B(t), uniquely defined by the following properties³:

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- (ii) B(t) is stationary, i.e., for $t > s \ge 0$ the increment B(t) B(s) has the same distribution as B(t s).
- (iii) B(t) has independent increments. That is, for all $t_n > t_{n-1} > \ldots > t_2 > t_1$, the random variables $B(t_n) B(t_{n-1}), \ldots, B(t_2) B(t_1), B(t_1)$ are independently distributed (i.e., their joint distribution factorizes).



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- (iv) B(t) has Gaussian distribution with variance t for all $t \in (0, \infty)$.
- (v) B(t) is continuous with probability 1.

The probability distribution \mathbb{P} governing the driving process B(t) is commonly known as the Wiener measure.

SDEs in physicist's notation

$$dX(t) = u dt + \sqrt{2D} dB(t). \tag{1.25}$$

Although the derivative $\xi(t) = dB/dt$ is not well-defined mathematically, Eq. (1.25) is in the physics literature often written in the form

$$\dot{X}(t) = u + \sqrt{2D}\,\xi(t). \tag{1.26}$$

The random driving function $\xi(t)$ is then referred to as Gaussian white noise, characterized by

$$\langle \xi(t) \rangle = 0 , \qquad \langle \xi(t)\xi(s) \rangle = \delta(t-s), \qquad (1.27)$$

with $\langle \cdot \rangle$ denoting an average with respect to the Wiener measure.

Stochastic differential calculus

$$dX(t) = u dt + \sqrt{2D} dB(t). \tag{1.25}$$

Ito's formula Note that property (iv) implies that $\mathbb{E}[dB^2] = dt$. This justifies the following heuristic derivation of Ito's formula for the differential change of some real-valued function F(x)

$$dF(X(t)) := F(X(t+dt)) - F(X(t))$$

$$= F'(X(t)) dX + \frac{1}{2}F''(X(t)) dX^{2} + \dots$$

$$= F'(X(t)) dX + \frac{1}{2}F''(X(t)) \left[u dt + \sqrt{2D} dB \right]^{2} + \dots$$

$$= F'(X(t)) dX + DF''(X(t)) dB^{2} + O(dt^{3/2}); \tag{1.28}$$

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hence, in a probabilistic sense, one has to leading order in dt

$$dF(X(t)) = F'(X(t)) dX + DF''(X(t)) dt$$

= $[uF'(X(t)) + DF''(X(t))] dt + F'(X(t)) \sqrt{2D} dB(t).$ (1.29)

It is crucial to note that, due to the choice of the expansion point, the coefficient F'(X) in front of dB(t) is to be evaluated at X(t). This convention is the so-called Ito integration rule. In particular, it is important to keep in mind that nonlinear transformations of Ito SDEs must feature second-order derivatives.

Numerical integration

Discretization dilemma To clarify the importance of discretization rules when dealing with SDEs, let us consider a simple generalization of Eq. (1.25), where drift u and diffusion coefficient D are position dependent. Adopting the Ito convention, the corresponding SDE reads

$$dX(t) = u(X) dt + \sqrt{2D(X)} * dB(t),$$
 (1.30a)

where from now on the *-symbol signals that D(X) is to be evaluated at X(t). The simplest numerical integration procedure for Eq. (1.30a) is the stochastic Euler scheme

$$X(t + dt) = X(t) + u(X(t)) dt + \sqrt{2D(X(t))} \sqrt{dt} Z(t), \qquad (1.30b)$$

$$Z(t) \sim \mathcal{N}(0,1)$$

When you see an equation like (1.30a), then always ask which discretization rule has been adopted!

Ito vs. backward-Ito

Compare

$$dX(t) = u(X) dt + \sqrt{2D(X)} * dB(t),$$
 (1.30a)

$$X(t + dt) = X(t) + u(X(t)) dt + \sqrt{2D(X(t))} \sqrt{dt} Z(t),$$
 (1.30b)

with

the so-called backward Ito SDE with coefficients u_B and D_B , denoted by

$$dX(t) = u_B(X) dt + \sqrt{2D_B(X)} \bullet dB(t), \qquad (1.31a)$$

is defined as the upper Riemann sum⁶

$$X(t+dt) = X(t) + u_B(X(t+dt)) dt + \sqrt{2D_B(X(t+dt))} \sqrt{dt} Z(t).$$
 (1.31b)

do NOT give same results when $dt \rightarrow 0$

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In particular

$$dF(X) = F'(X) \bullet dX - D_B F''(X) dt$$

$$= [u_B F'(X) - D_B F''(X)] dt + F'(X) \sqrt{2D_B} \bullet dB(t).$$
(1.32)

Stratonovich SDE

Another discretization convention, that is popular in the physics literature is the Stratonovich-Fisk discretization, denoted by

$$dX(t) = u_S(X) dt + \sqrt{2D_S(X)} \circ dB(t), \qquad (1.34a)$$

and defined as the mean value of lower and upper Riemann sum⁸

$$X(t+dt) = X(t) + \frac{u_S(X(t)) + u_S(X(t+dt))}{2} dt + \frac{\sqrt{2D_S(X(t))} + \sqrt{2D_S(X(t+dt))}}{2} \sqrt{dt} Z(t).$$
 (1.34b)

From a numerical perspective, the non-anticipatory Ito scheme (1.30b) is advantageous for it allows to compute the new position directly from the previous one. For analytical calculations, the Stratonovich-Fisk scheme is somewhat preferable as it preserves the rules of ordinary differential calculus,⁹

$$dF(X) = F'(X) \circ dX(t) \tag{1.36}$$

each SDE formulation has advantages & disadvantages

Summary

Ito

$$dX(t) = u(X) dt + \sqrt{2D(X)} * dB(t),$$
 (1.30a)

$$X(t + dt) = X(t) + u(X(t)) dt + \sqrt{2D(X(t))} \sqrt{dt} Z(t), \tag{1.30b}$$

backward-lto

$$dX(t) = u_B(X) dt + \sqrt{2D_B(X)} \bullet dB(t)$$

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Stratonovich

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Since other types of SDEs can be transformed into an equivalent Ito SDE, we shall focus in this part on discussing how one can derive a Fokker-Planck equation (FPE) for the probability density function (PDF) p(t, x) for a process X(t) described by the Ito SDE

$$dX(t) = u(X) dt + \sqrt{2D(X)} * dB(t).$$
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The PDF can be formally defined by

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$$\mathbb{E}[d[\delta(X-x)]] = \mathbb{E}[(\partial_X \delta(X-x)) dX + D(X) \partial_X^2 \delta(X(t)-x) dt]$$
$$= \mathbb{E}[(\partial_X \delta(X-x)) u(X) + D(X) \partial_X^2 \delta(X(t)-x)] dt.$$

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Furthermore, by recalling that

$$\partial_X \delta(X - x) = -\partial_x \delta(X - x), \tag{1.40}$$

we may write

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Using another property of the δ -function

$$f(y)\delta(y-x) = f(x)\delta(y-x) \tag{1.41}$$

we obtain

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Combining this with Eq. (1.39) yields the Fokker-Planck (or Smoluchowski) equation

$$\partial_t p = -\partial_x \left\{ u(x) \, p - \partial_x [D(x)p] \right\}. \tag{1.42}$$

Ito-FPE

$$\partial_t p = -\partial_x \left\{ u(x) \, p - \partial_x [D(x)p] \right\}$$

Backward-Ito FPE

For comparison, an analogous calculation for the backward-Ito SDE

$$dX(t) = u_B(X) dt + \sqrt{2D_B(X)} \bullet dB(t), \qquad (1.43)$$

gives

$$\partial_t p = -\partial_x \left[u_B(x) \, p - D_B(x) \, \partial_x p \right]. \tag{1.44}$$

Compared with the Ito FPE (1.42), the diffusion coefficient D_B now enters in front of the gradient $\partial_x p$. Note, however, that the two different FPEs coincide if one identifies the coefficients as in Eq. (1.33):

$$u = u_B + \partial_x D_B, \qquad D = D_B$$