

Topological Defects

18.S995 - L23

Order Parameters, Broken Symmetry, and Topology

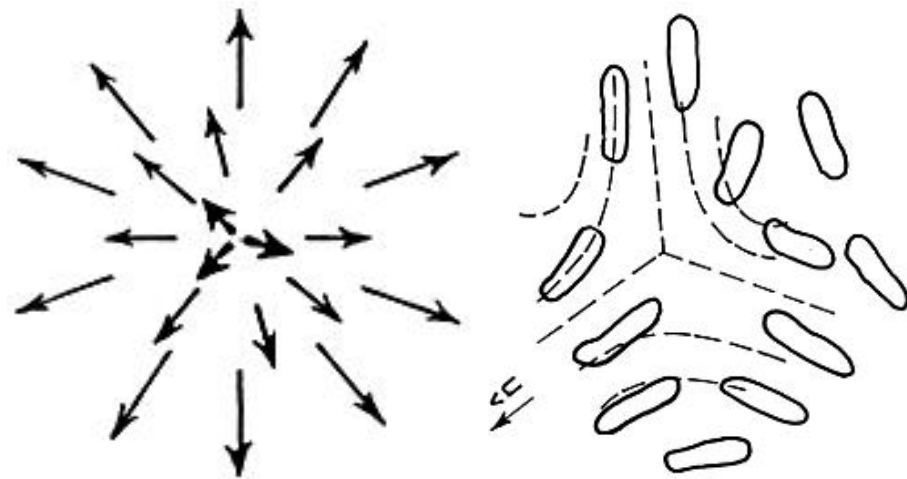
James P. Sethna

*Laboratory of Applied Physics, Technical University of Denmark,
DK-2800 Lyngby, DENMARK, and NORDITA, DK-2100 Copenhagen Ø,
DENMARK and Laboratory of Atomic and Solid State Physics (LASSP),
Clark Hall, Cornell University, Ithaca, NY 14853-2501, USA*

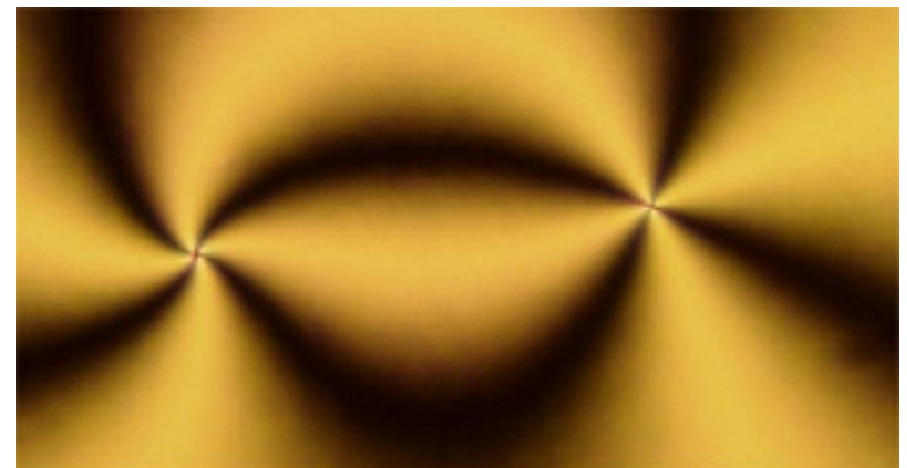
(Dated: May 27, 2003, 10:27 pm)

dunkel@mit.edu

Topological defects are discontinuities in order-parameter fields



- optical effects
- work hardening, etc



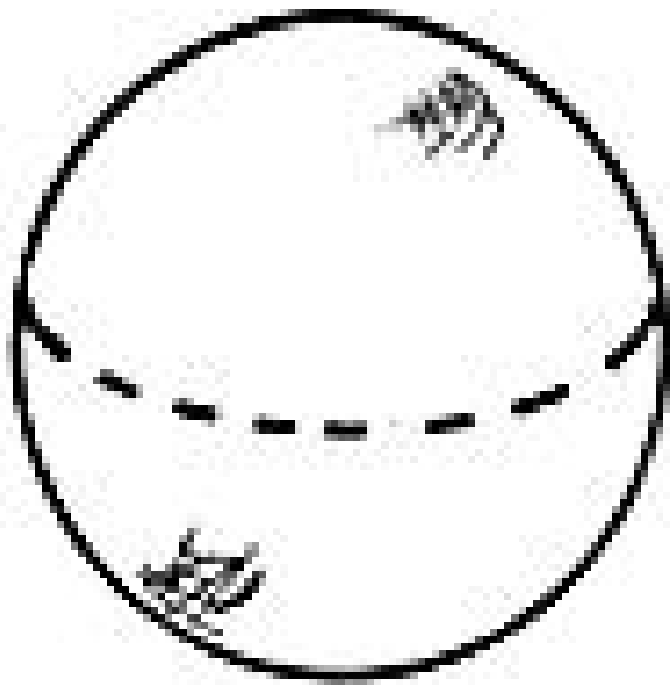
"umbilic defects" in a nematic liquid crystal

order = symmetry = invariance

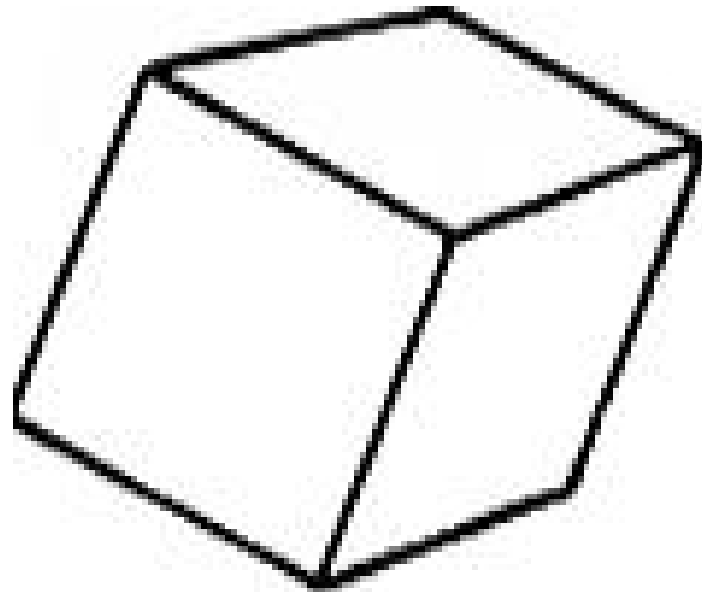
(under certain group actions)

symmetry groups can be discrete,
continuous, Lie-groups,

More or less symmetric ?



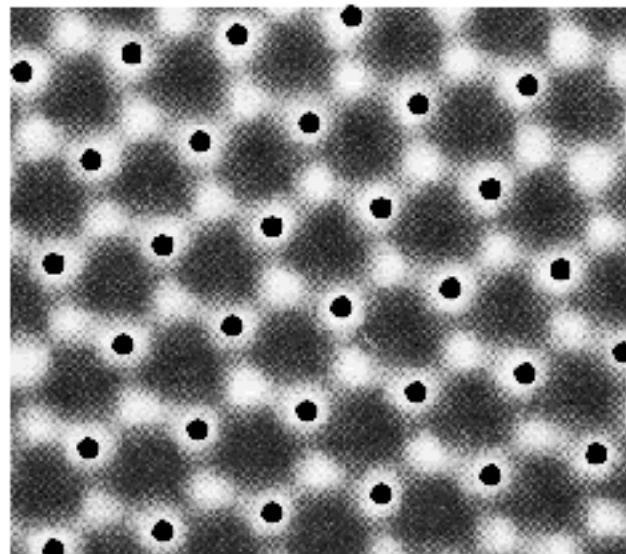
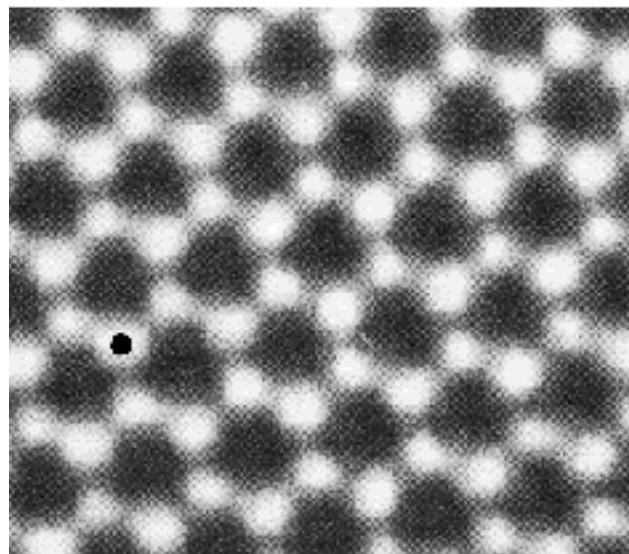
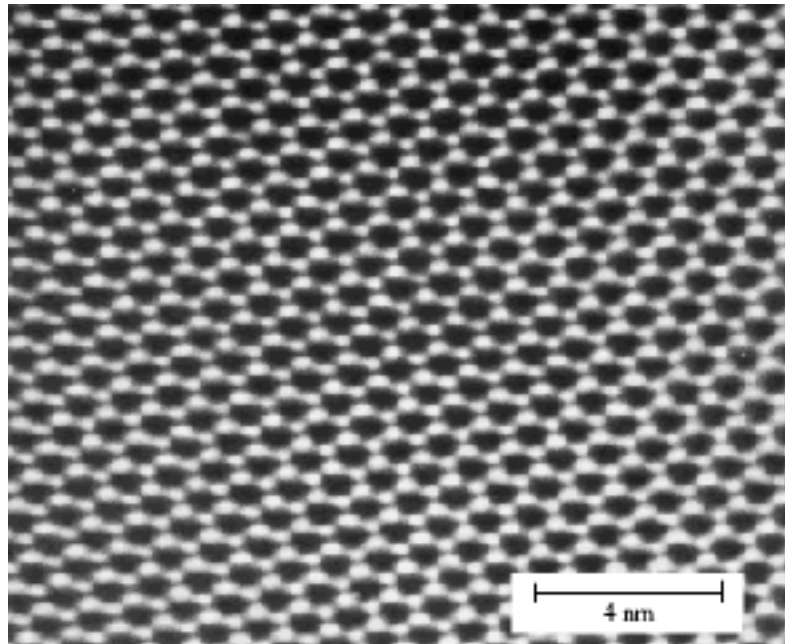
A



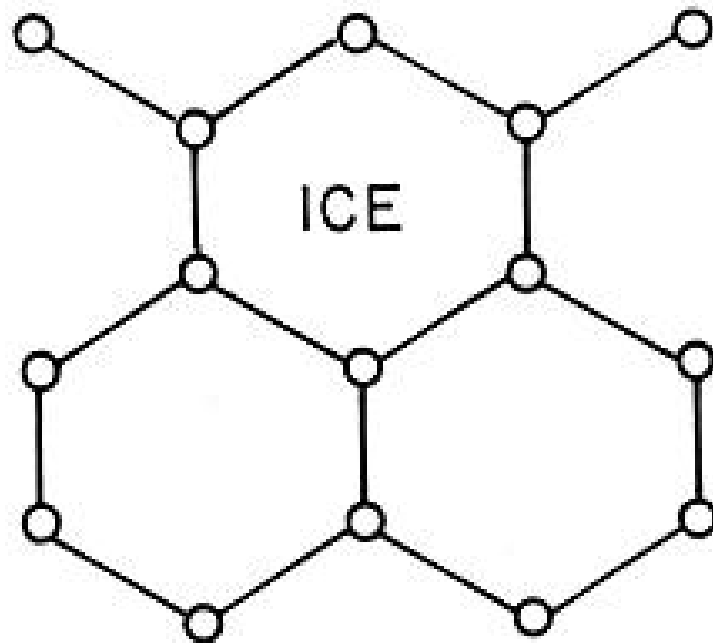
B

More or less symmetric ?

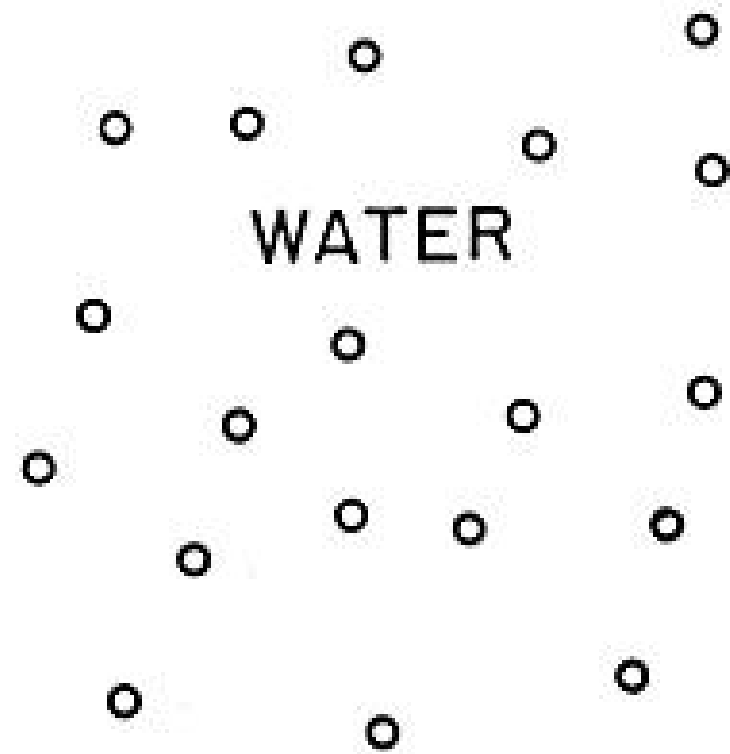
$\text{Mg}_2\text{Al}_4\text{Si}_5\text{O}_{18}$



More or less symmetric ?



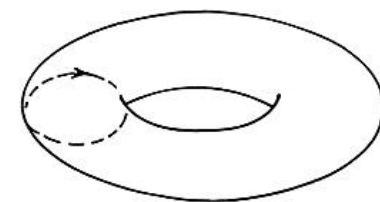
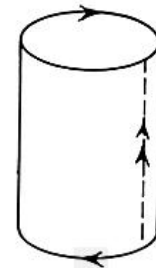
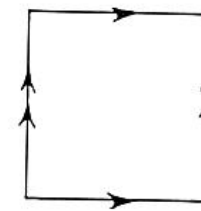
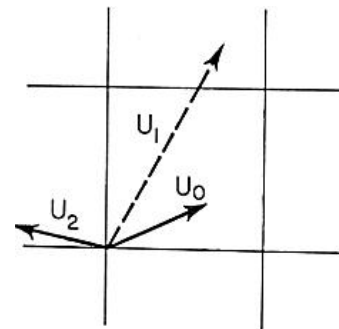
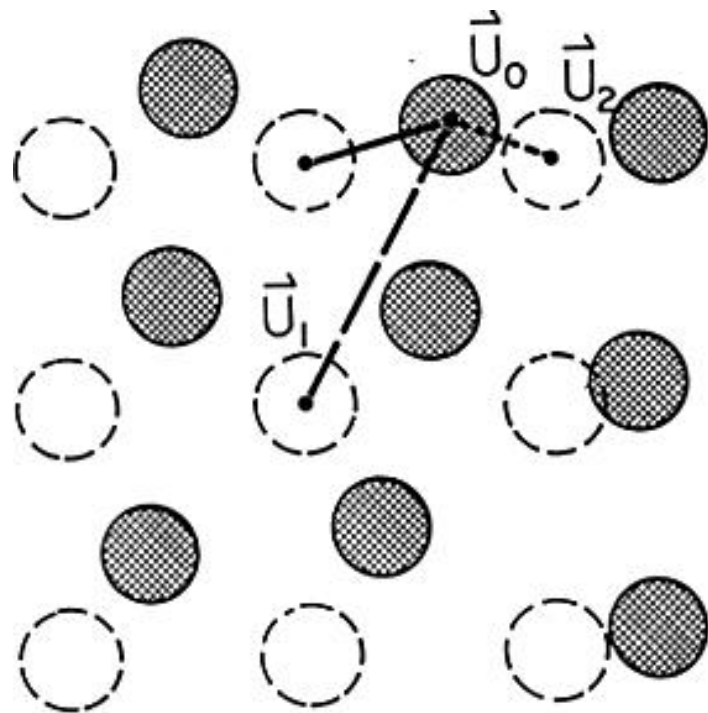
A



B

broken continuous
translation/rotation
symmetry (invariance)

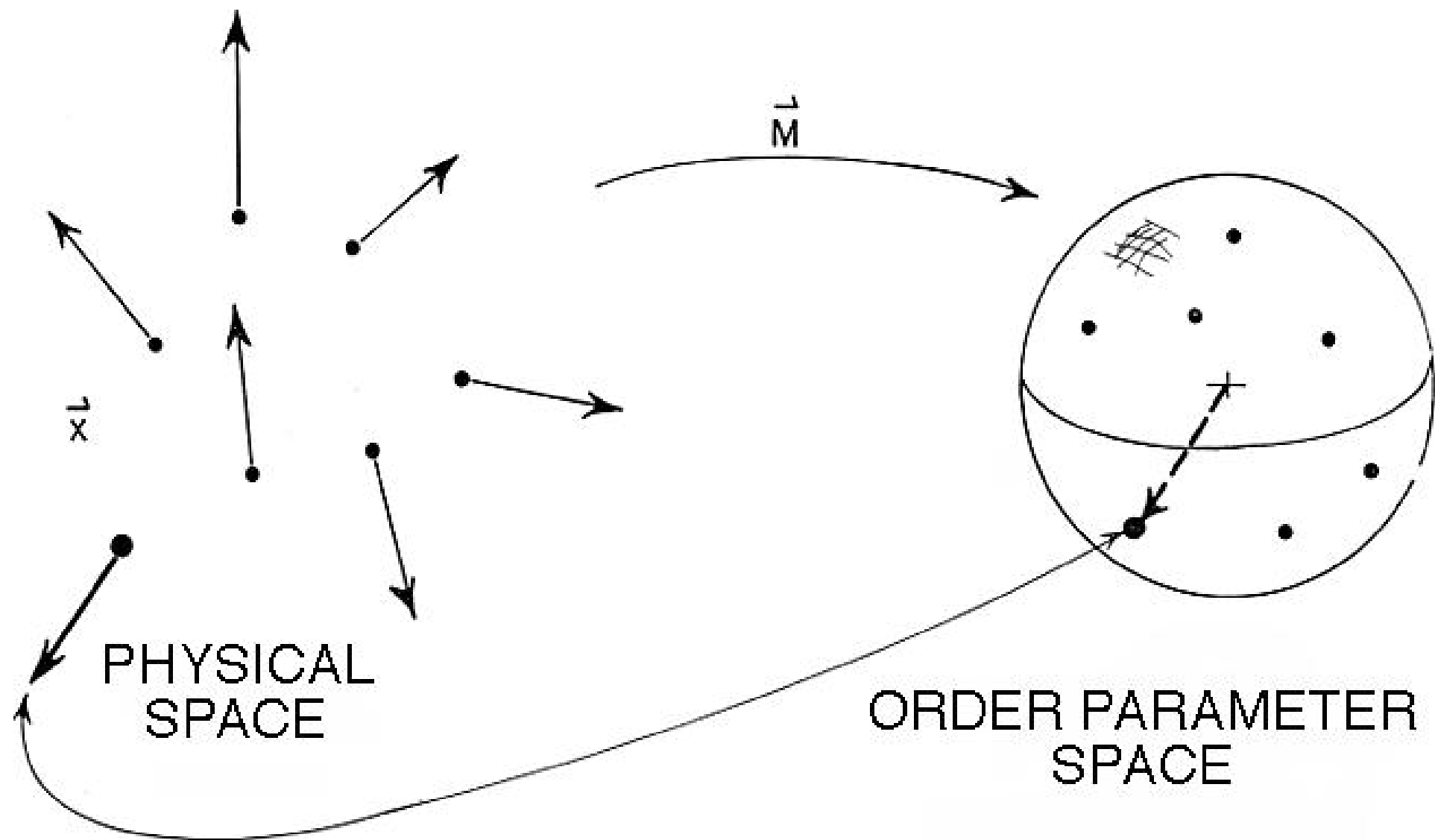
Order parameters: 2D crystal



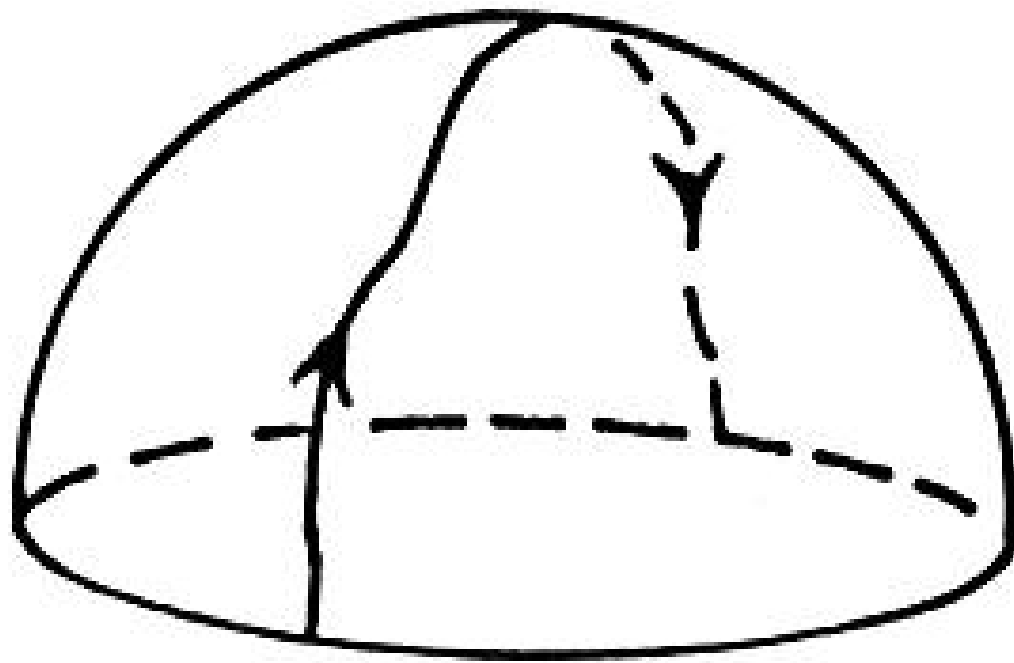
$$\vec{u} \equiv \vec{u} + a\hat{x} = \vec{u} + ma\hat{x} + na\hat{y}.$$

$$\mathcal{E} = \int dx (\kappa/2)(du/dx)^2.$$

Order parameters: magnets

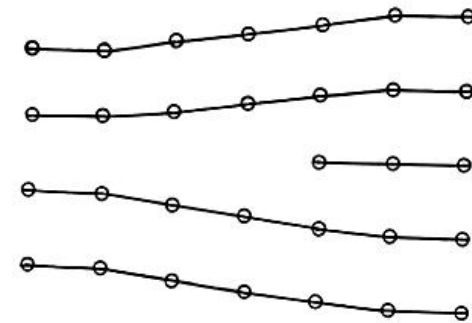
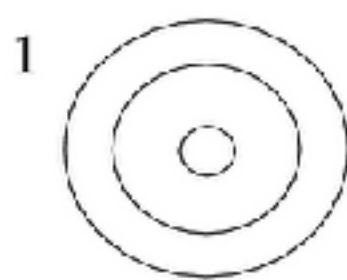
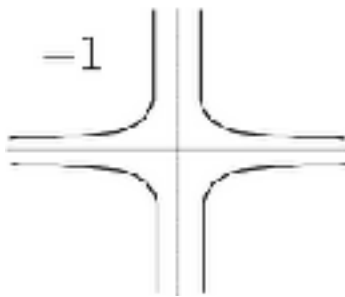
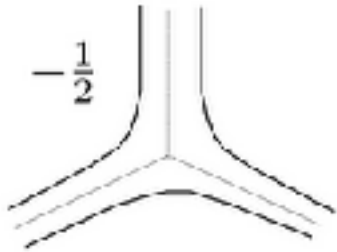
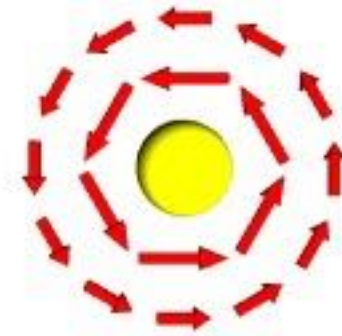
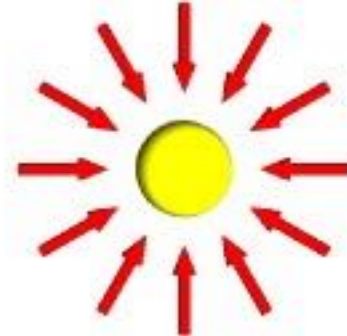
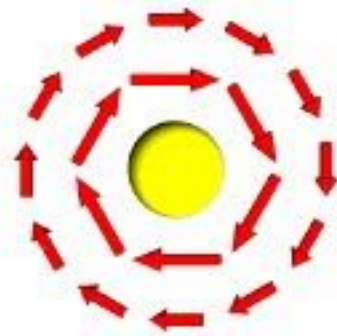
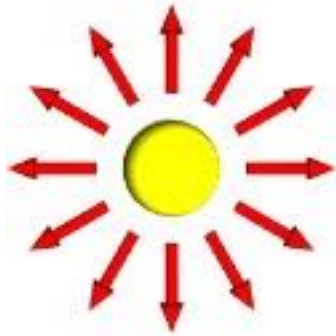


Order parameters: nematic liquid crystals

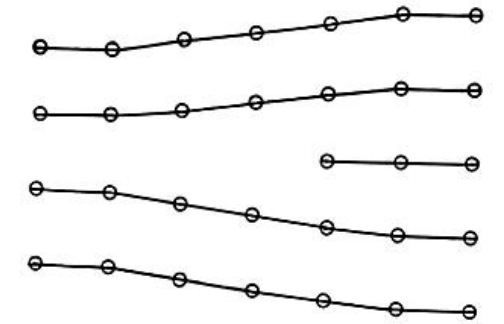


“projective plane” =
half-sphere
with opposite points on
equator identified

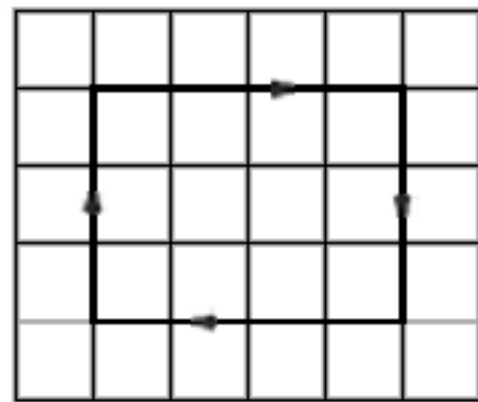
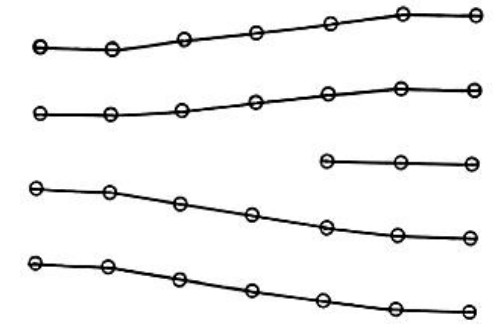
Topological defects



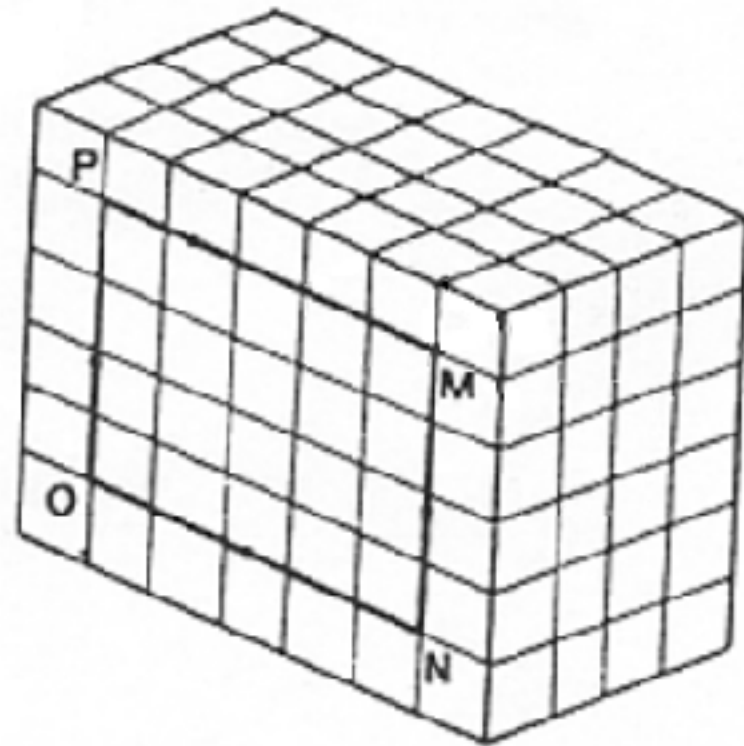
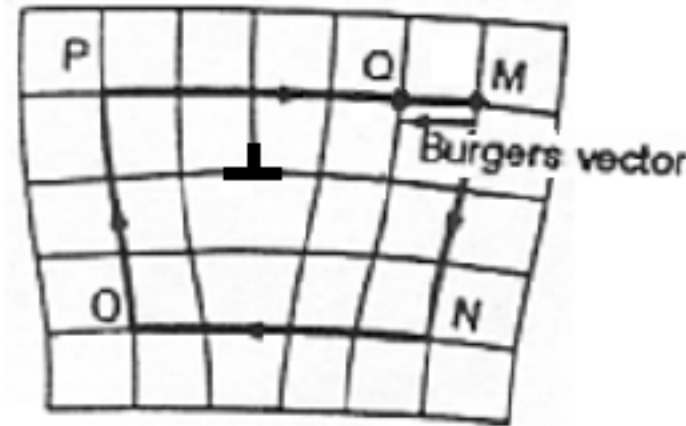
Work hardening



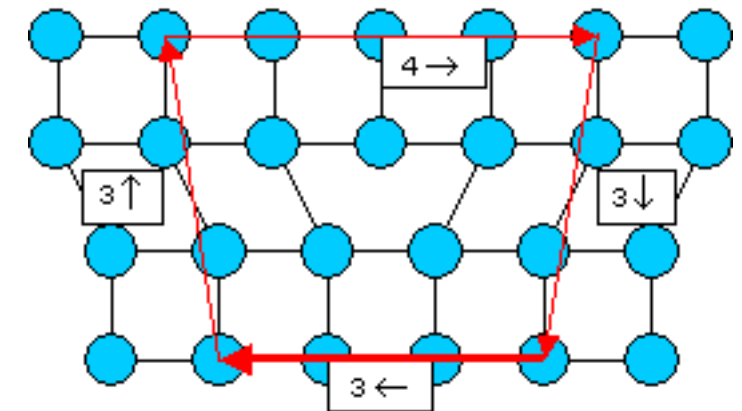
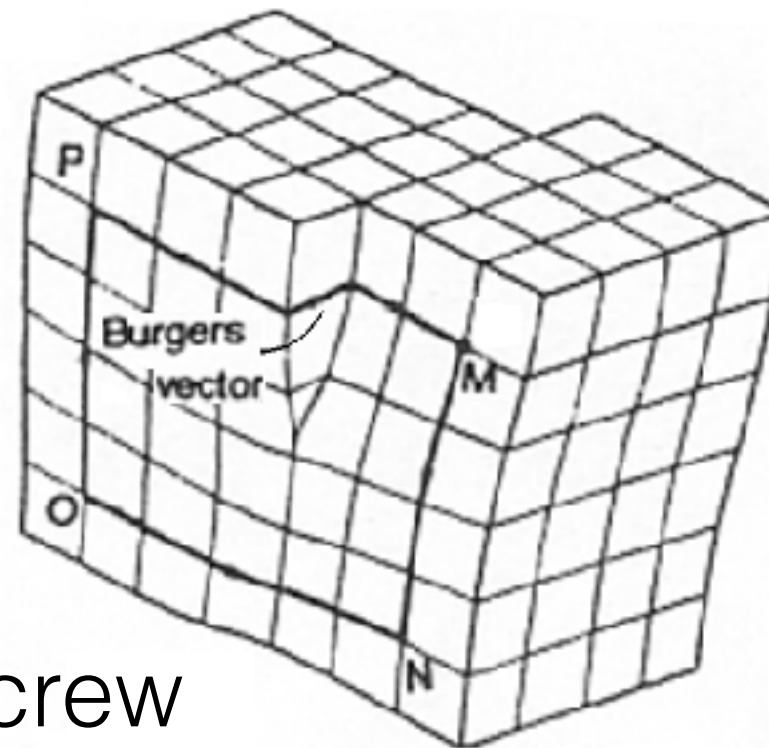
Dislocations



edge

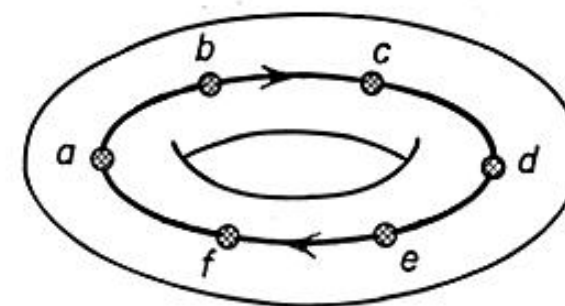
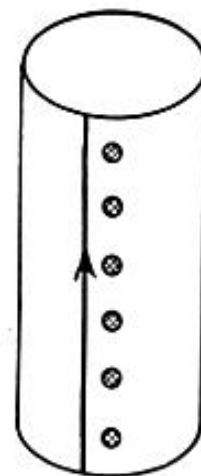
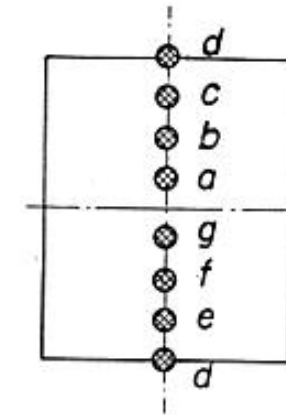
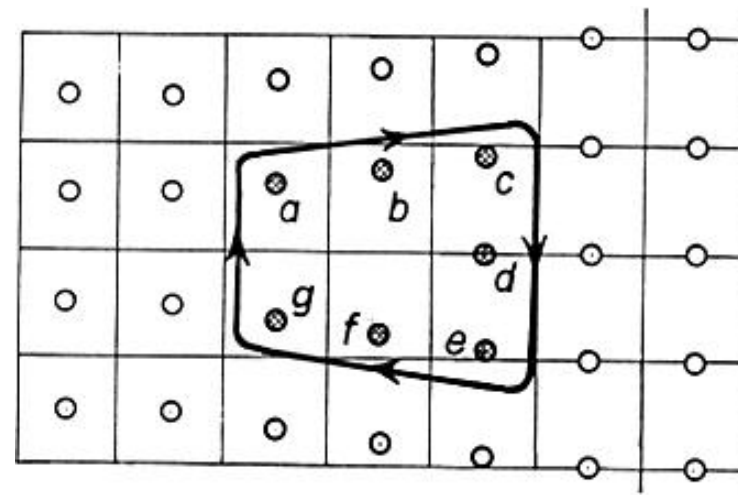
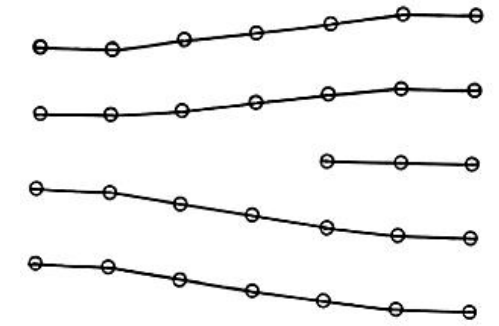


screw

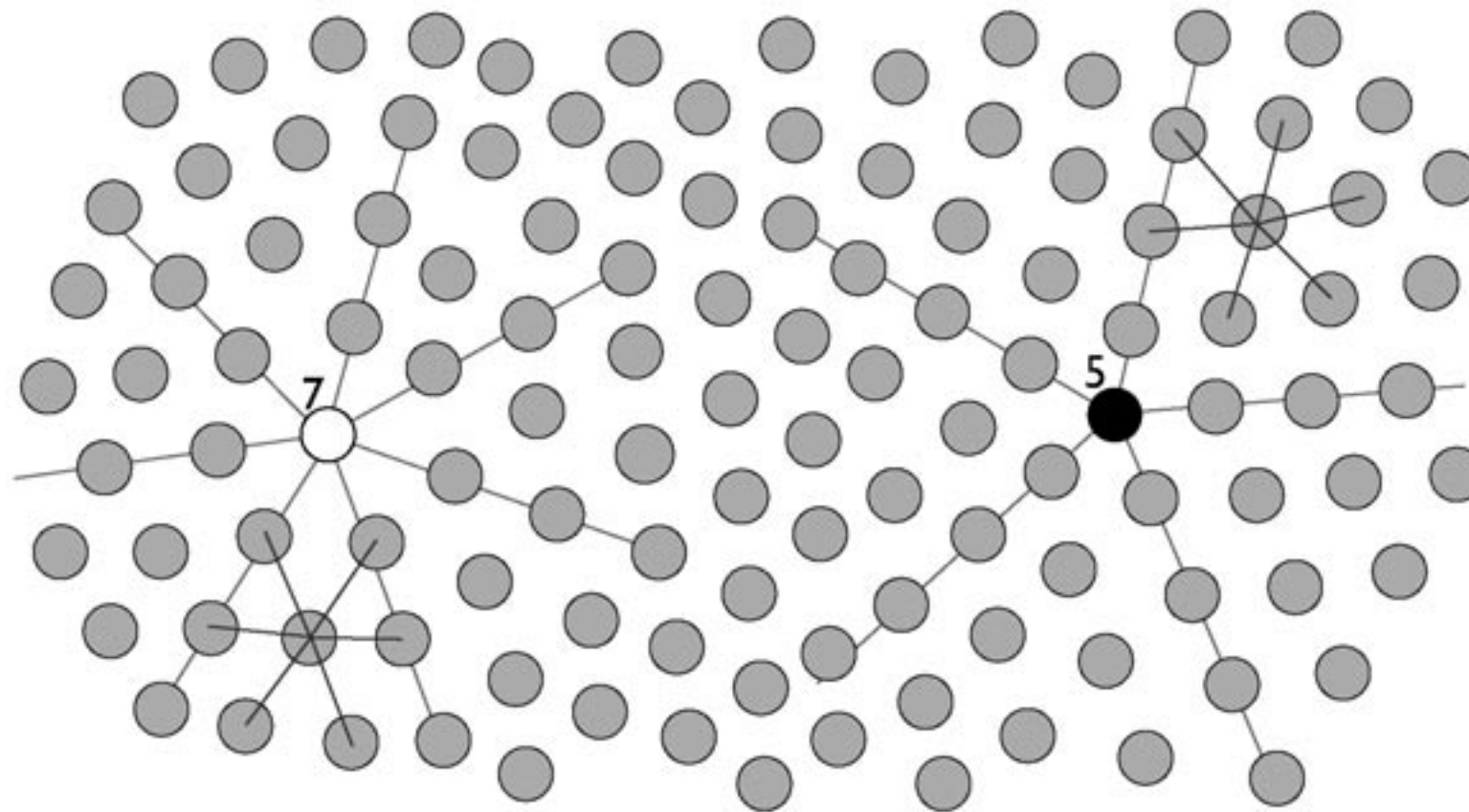


$$\|\mathbf{b}\| = (a/2)\sqrt{h^2 + k^2 + l^2}$$

Dislocations

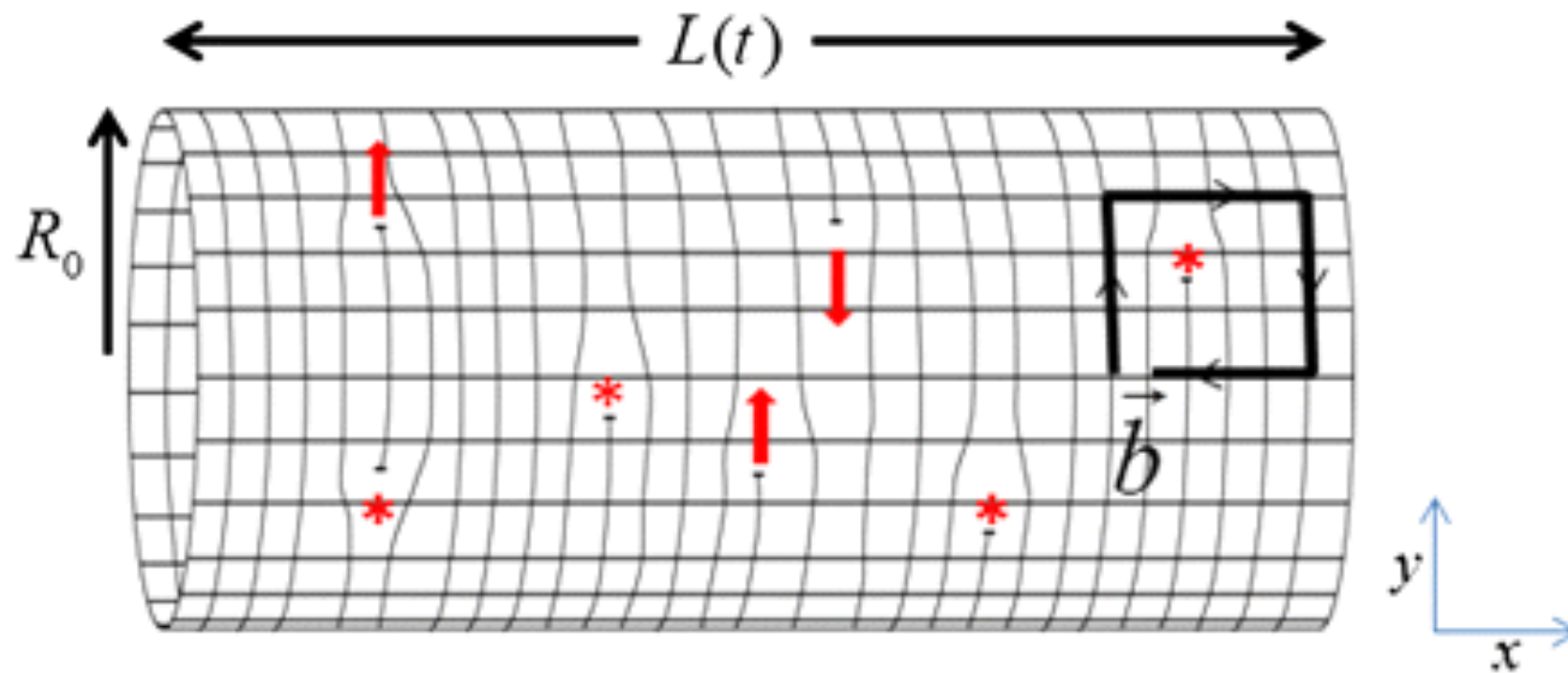


Disclination pair



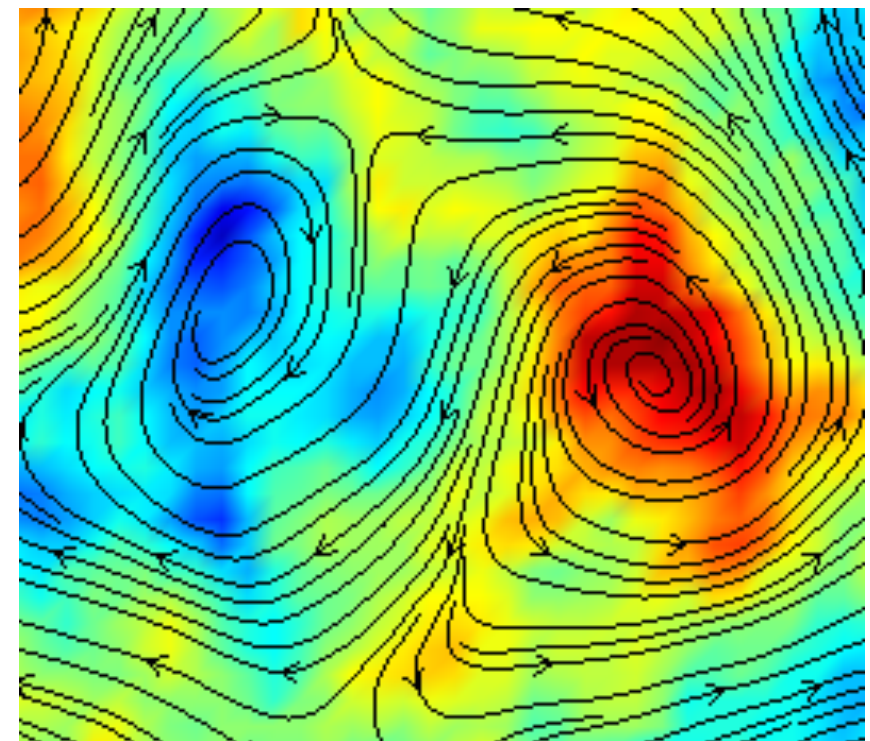
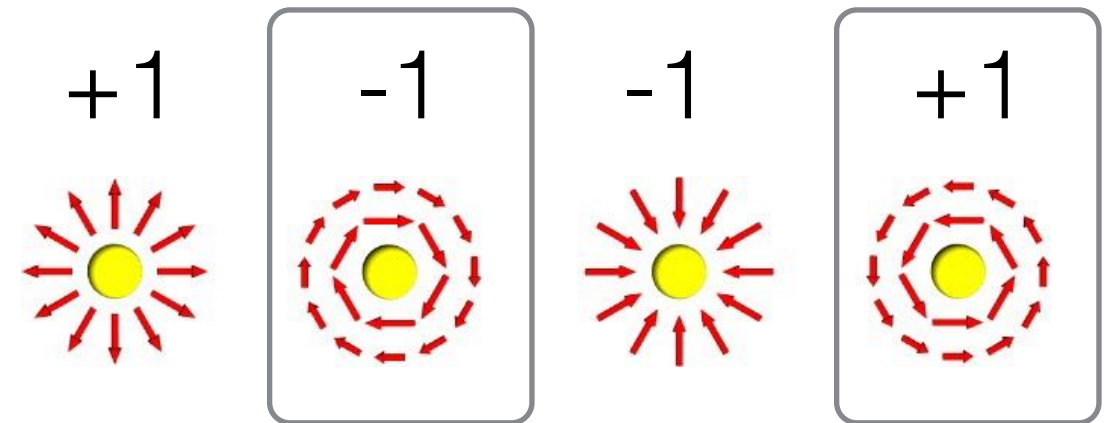
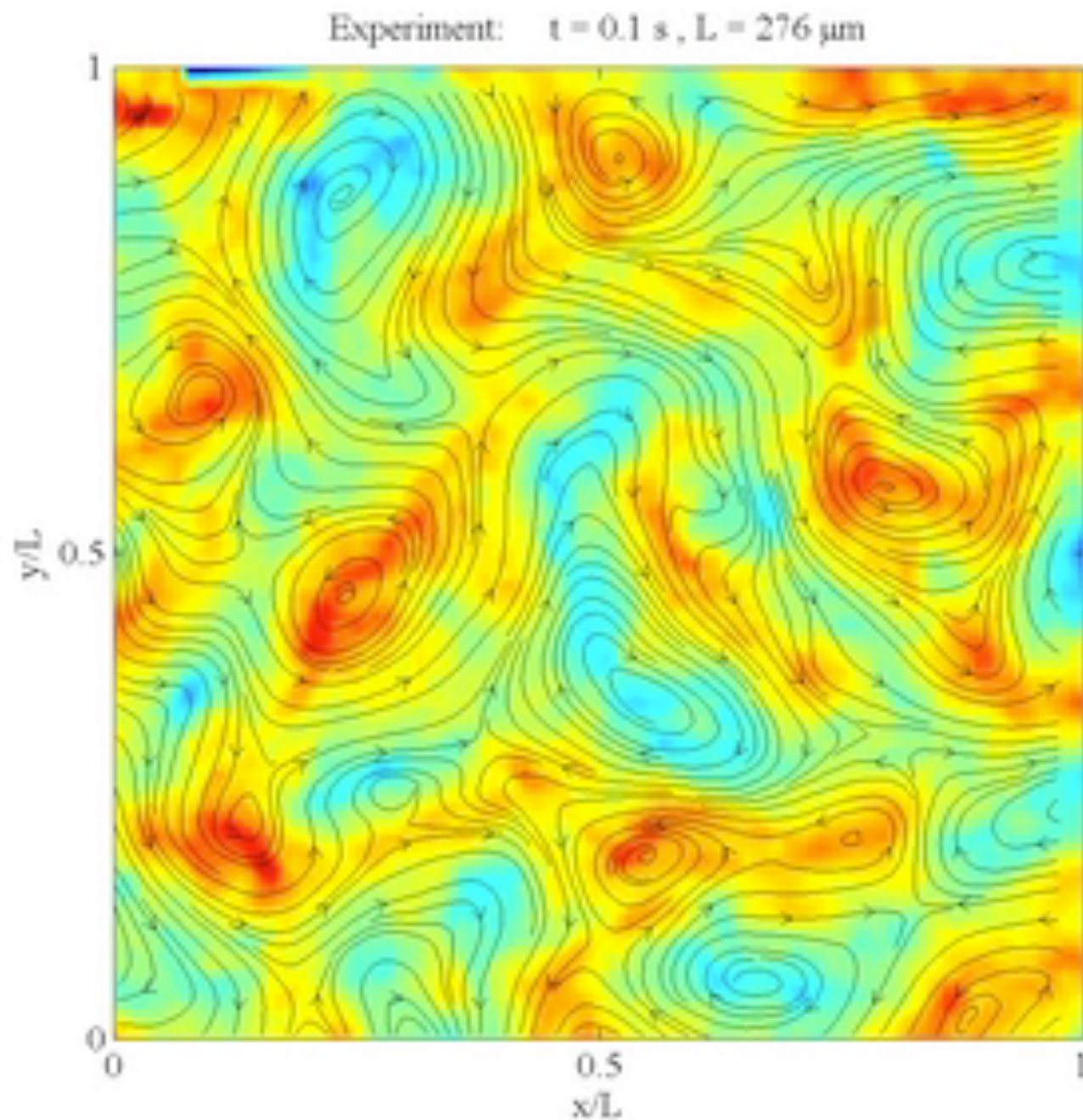
Dislocation-mediated growth of bacterial cell walls

Ariel Amir and David R. Nelson¹



Bacterial vortices

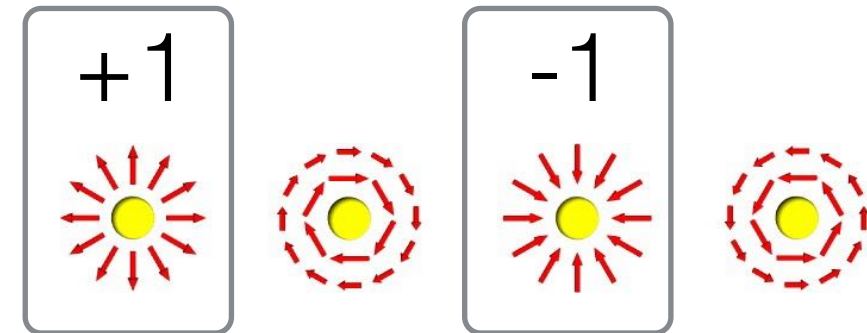
PIV



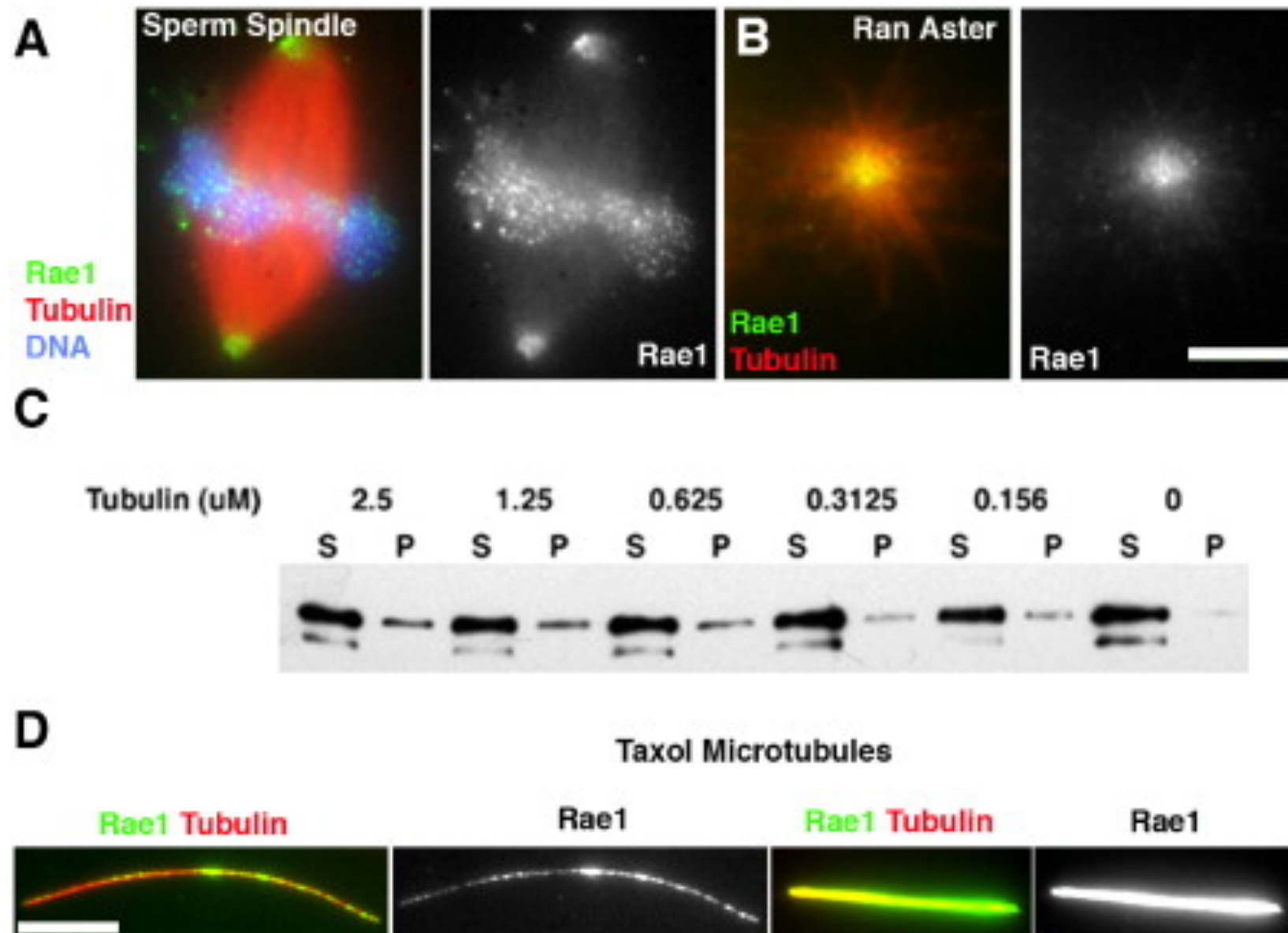
Dunkel et al PRL 2013

Microtubule asters

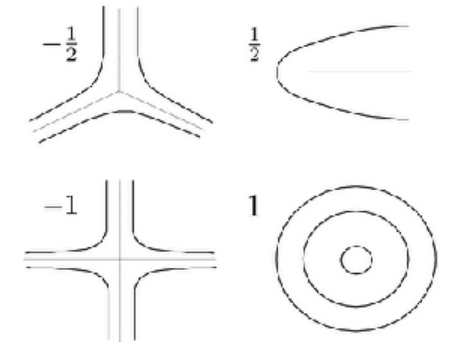
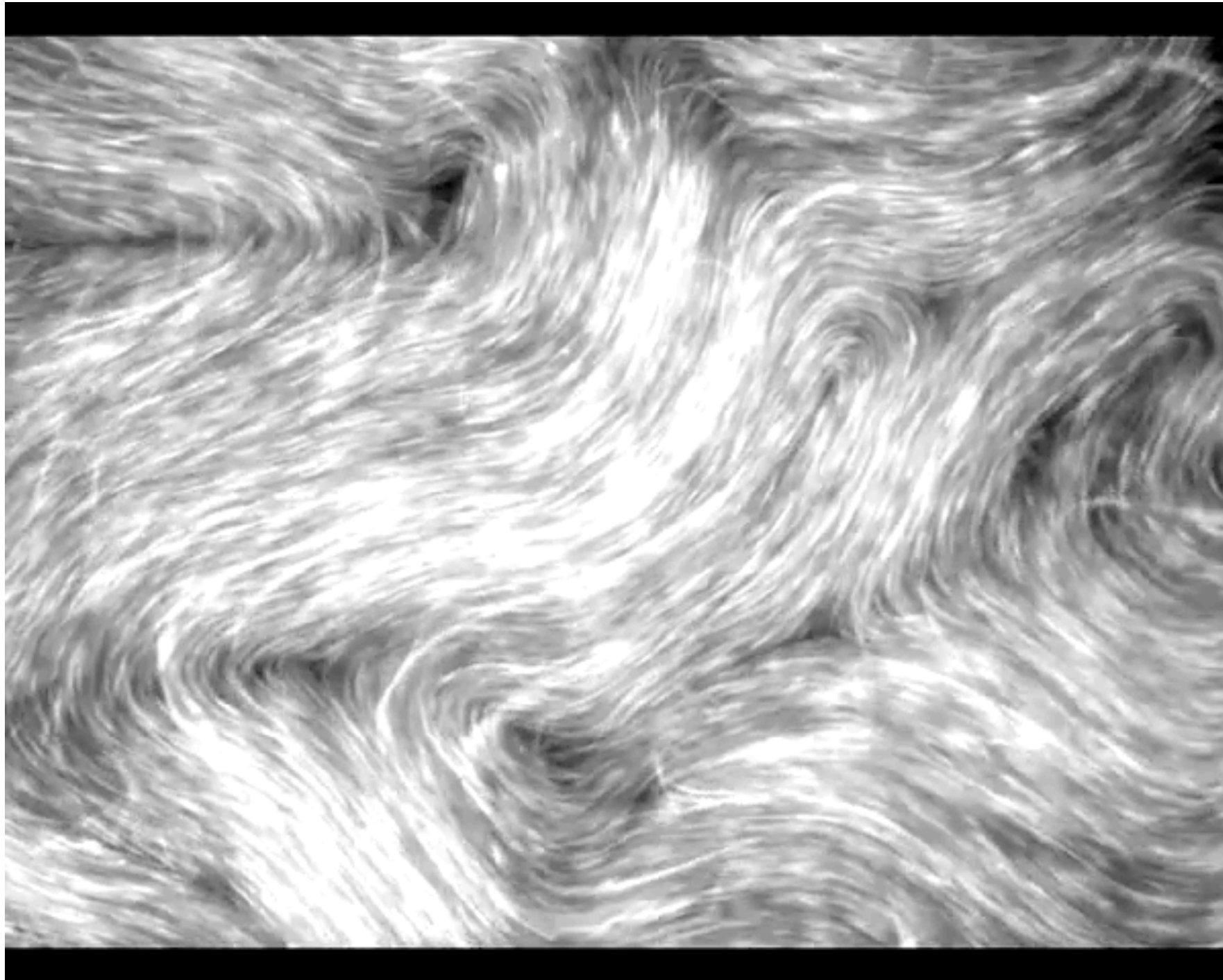
mitotic spindle organization



Blower et al (2005) Cell

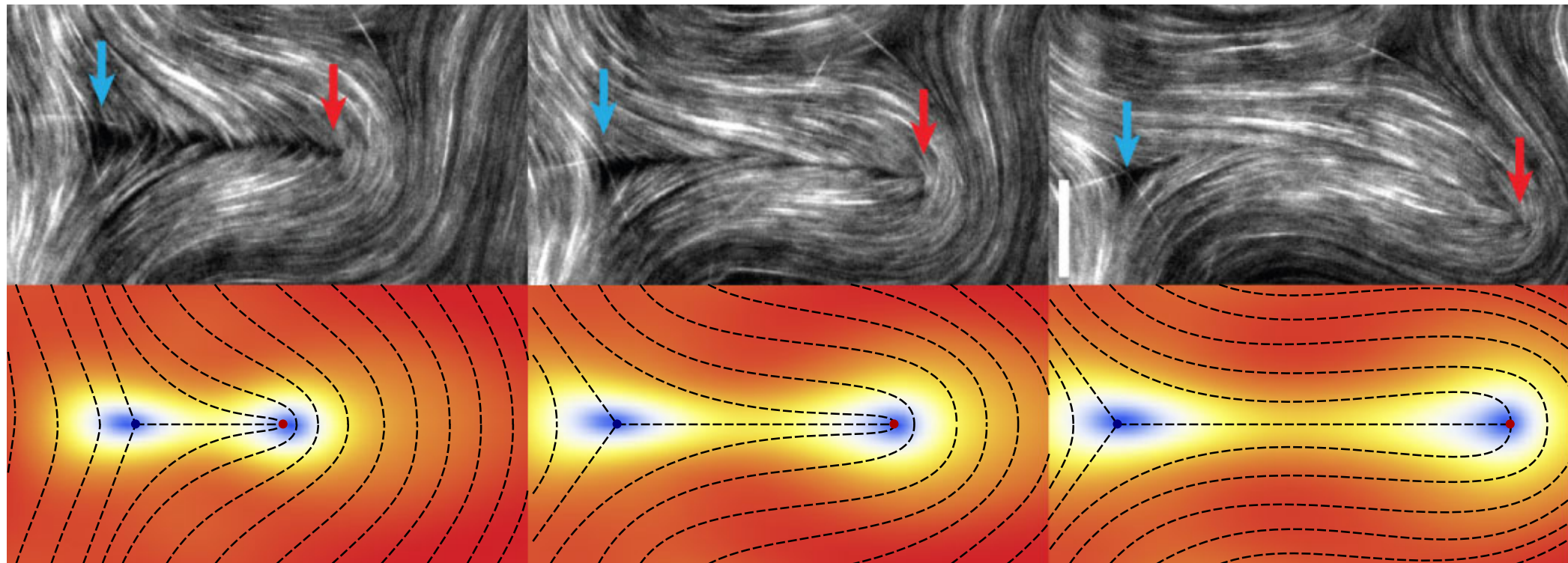
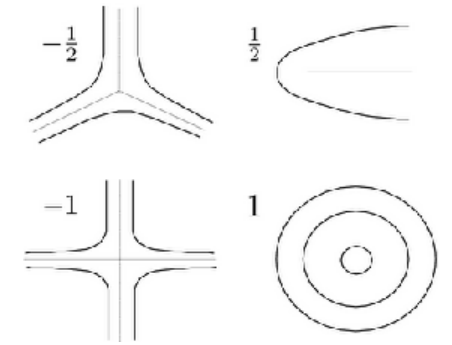


Active nematics



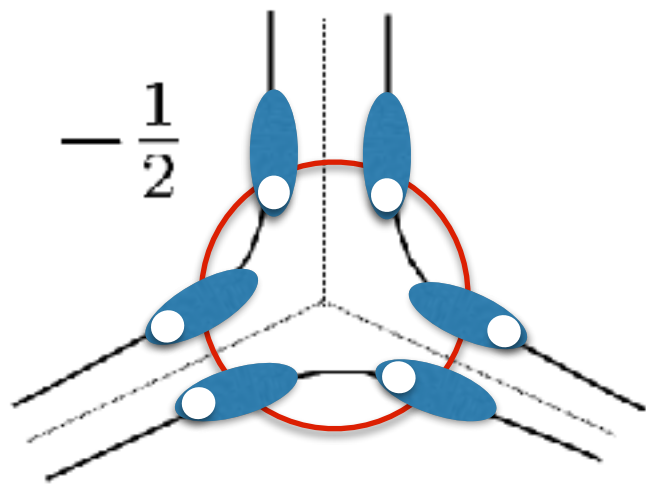
Dogic lab (Brandeis) Nature 2012

Active nematics

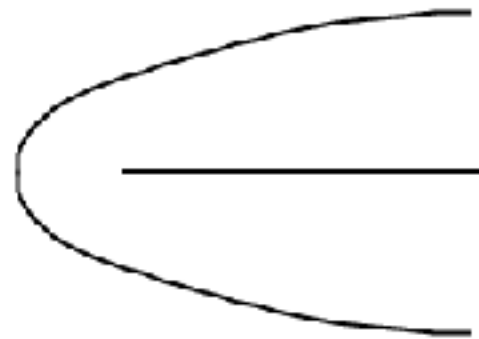


Giomi et al PRL 2012

Defects in nematics

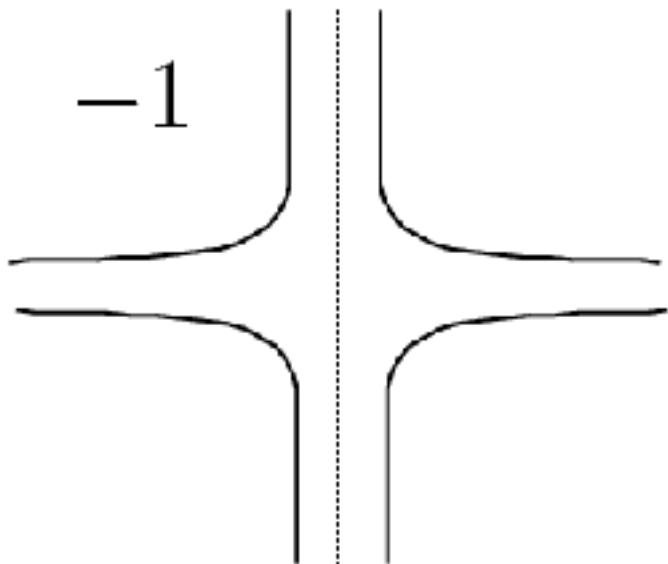


$\frac{1}{2}$

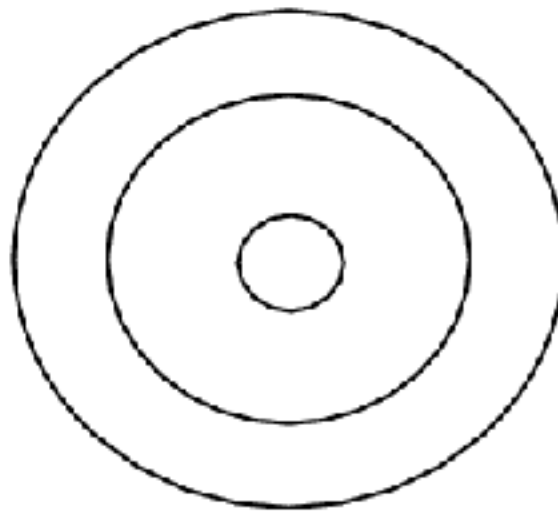


winding
number

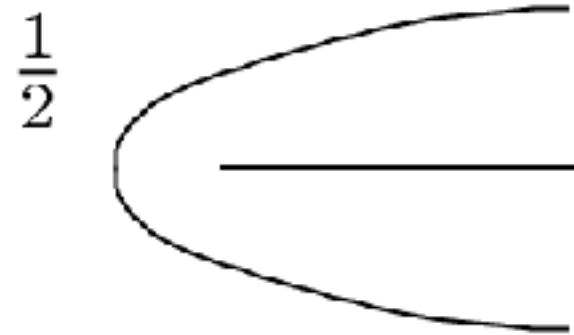
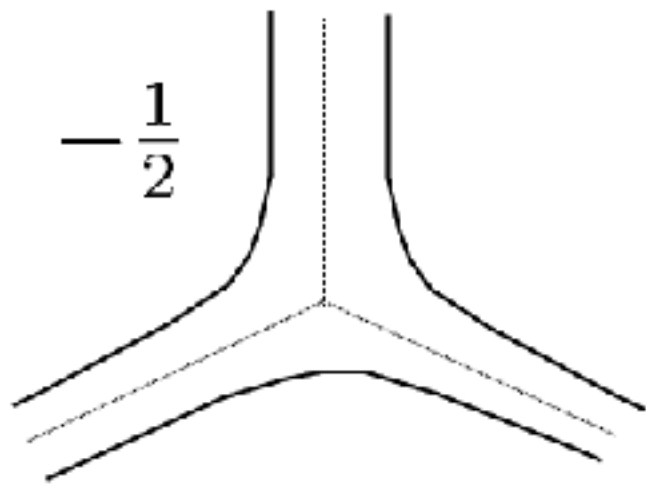
-1



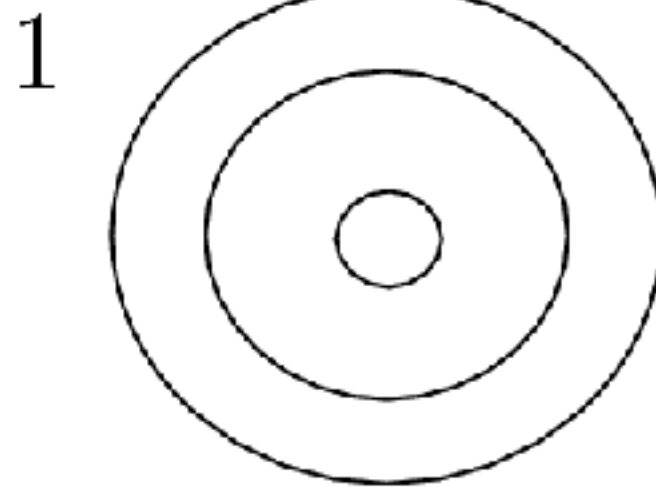
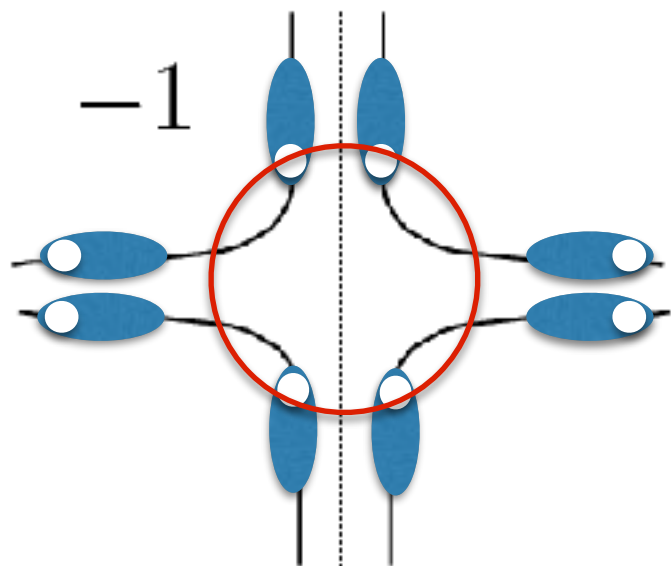
1



Defects in nematics



winding
number

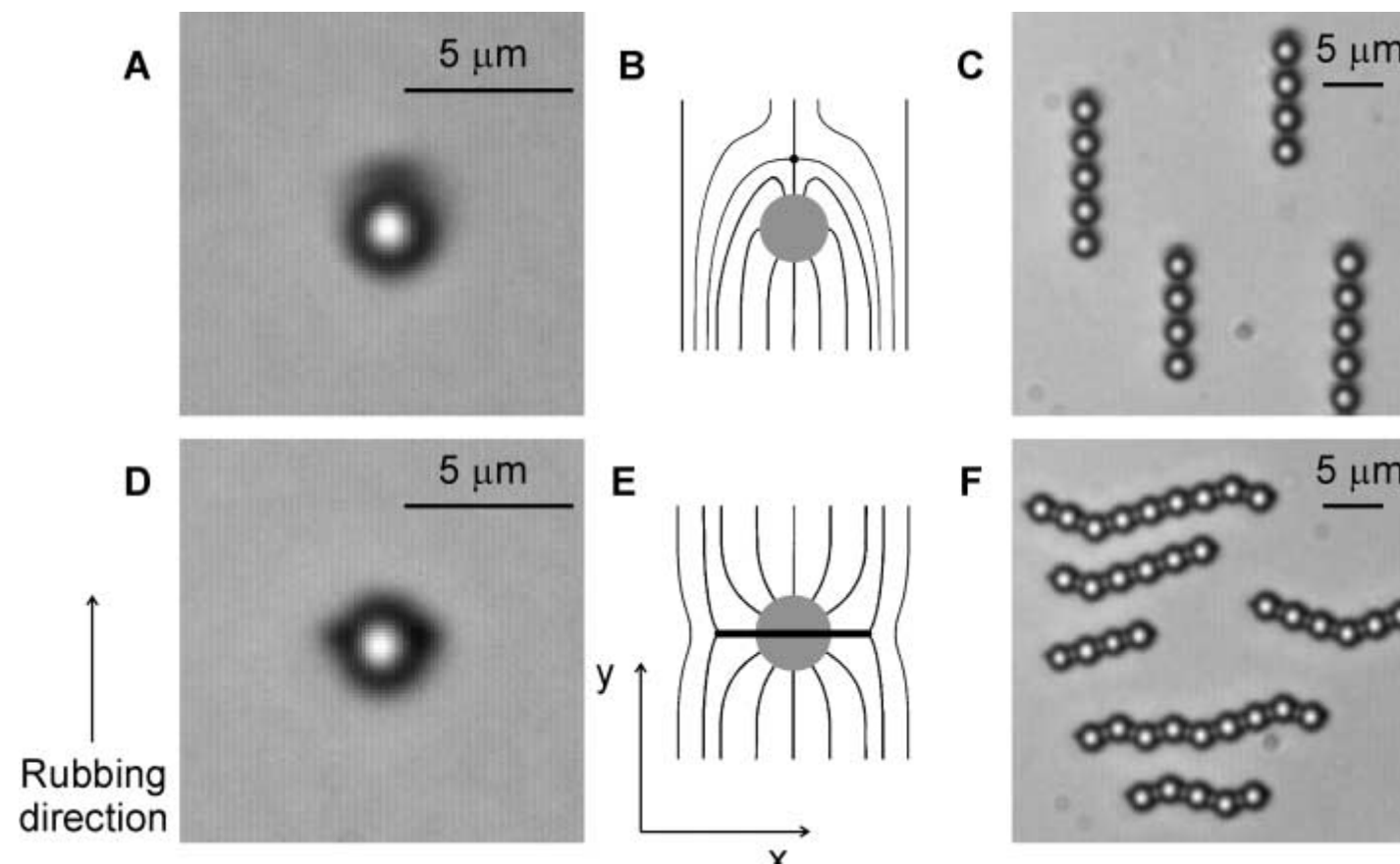


Two-Dimensional Nematic Colloidal Crystals Self-Assembled by Topological Defects

Igor Musevic *et al.*

Science **313**, 954 (2006);

DOI: 10.1126/science.1129660

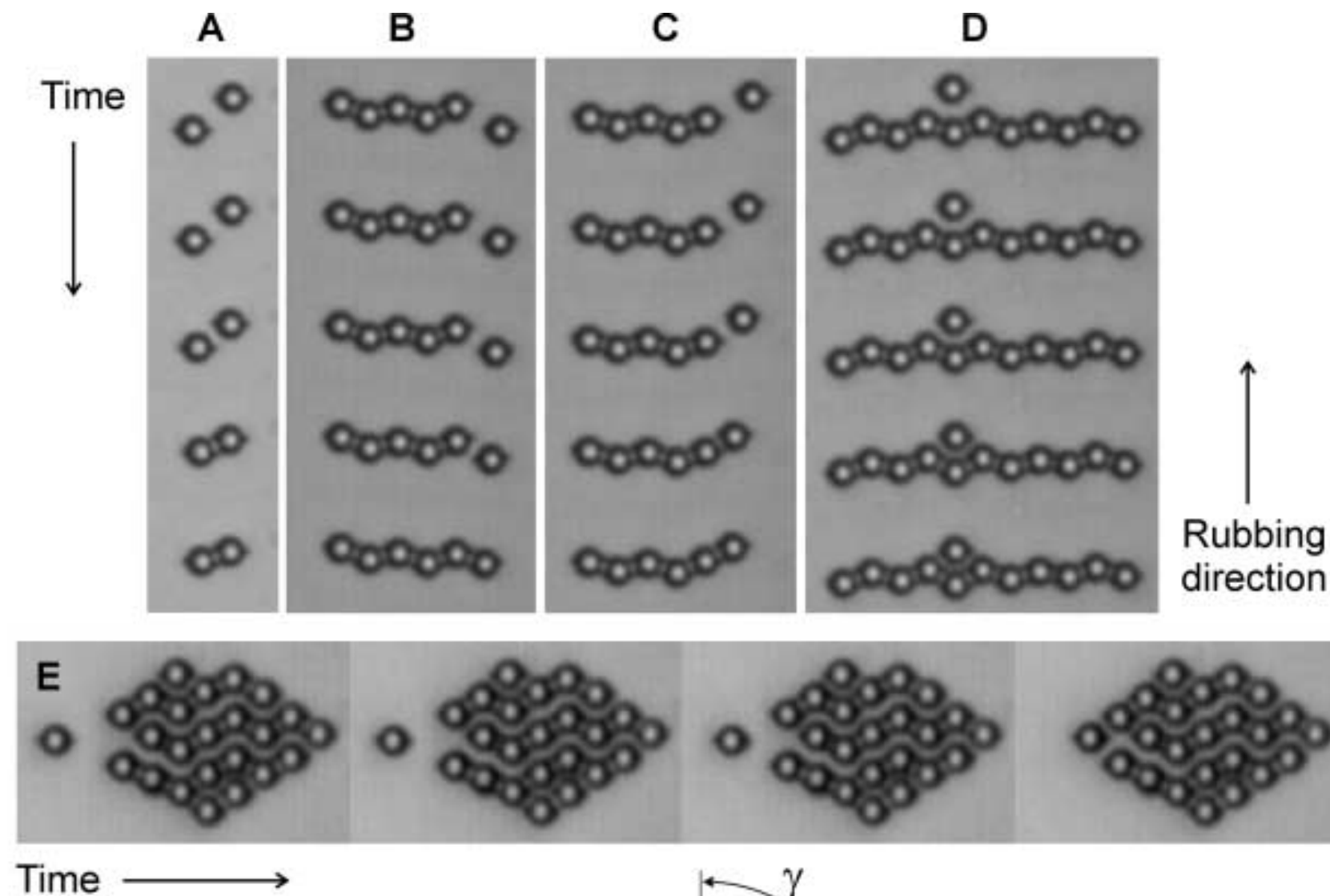


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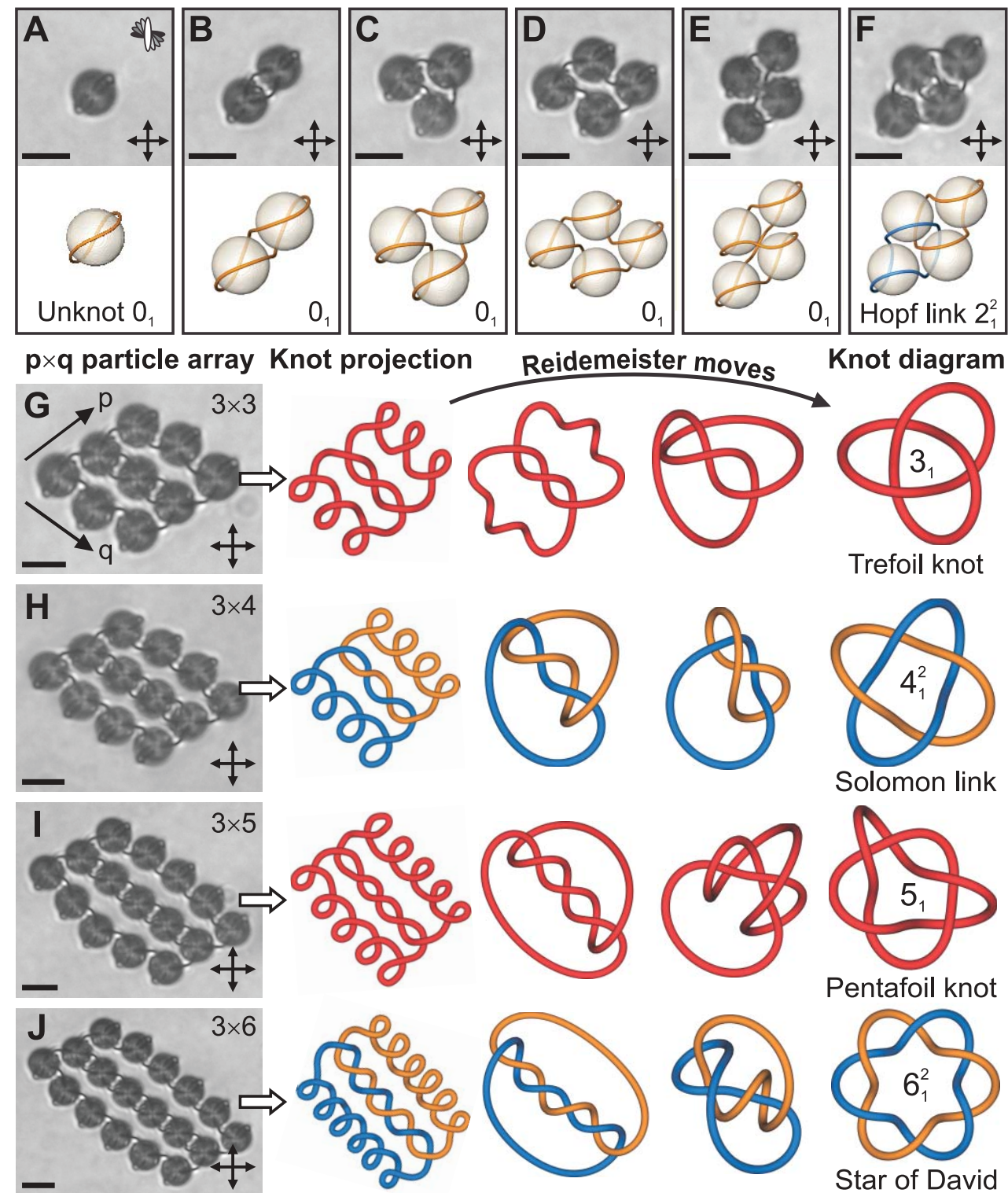


Reconfigurable Knots and Links in Chiral Nematic Colloids

Uros Tkalec *et al.*

Science **333**, 62 (2011);

DOI: 10.1126/science.1205705



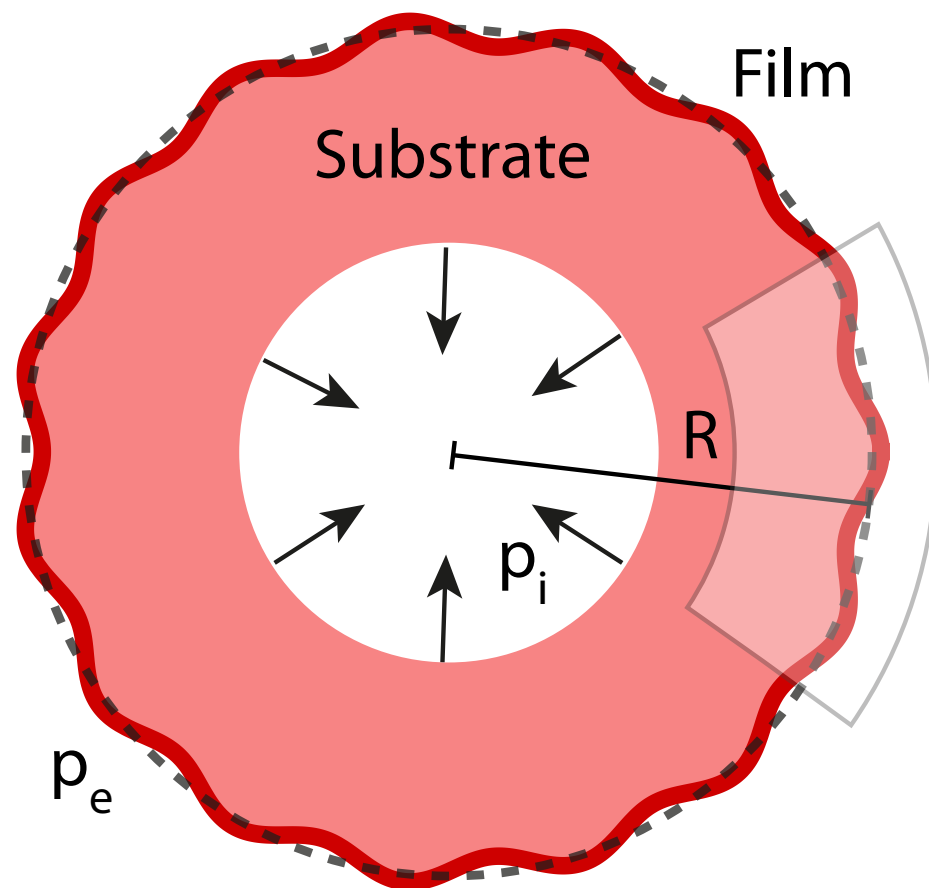
Experiment



Denis Terwagne



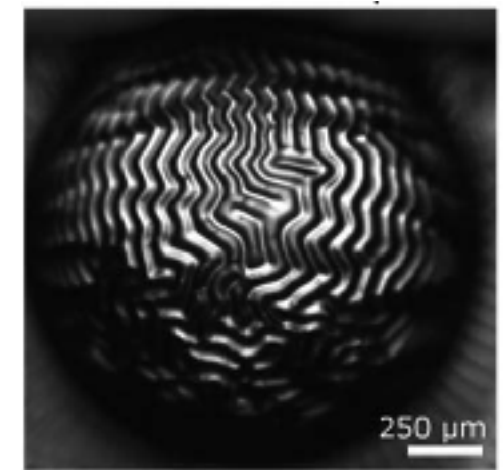
Pedro Reis



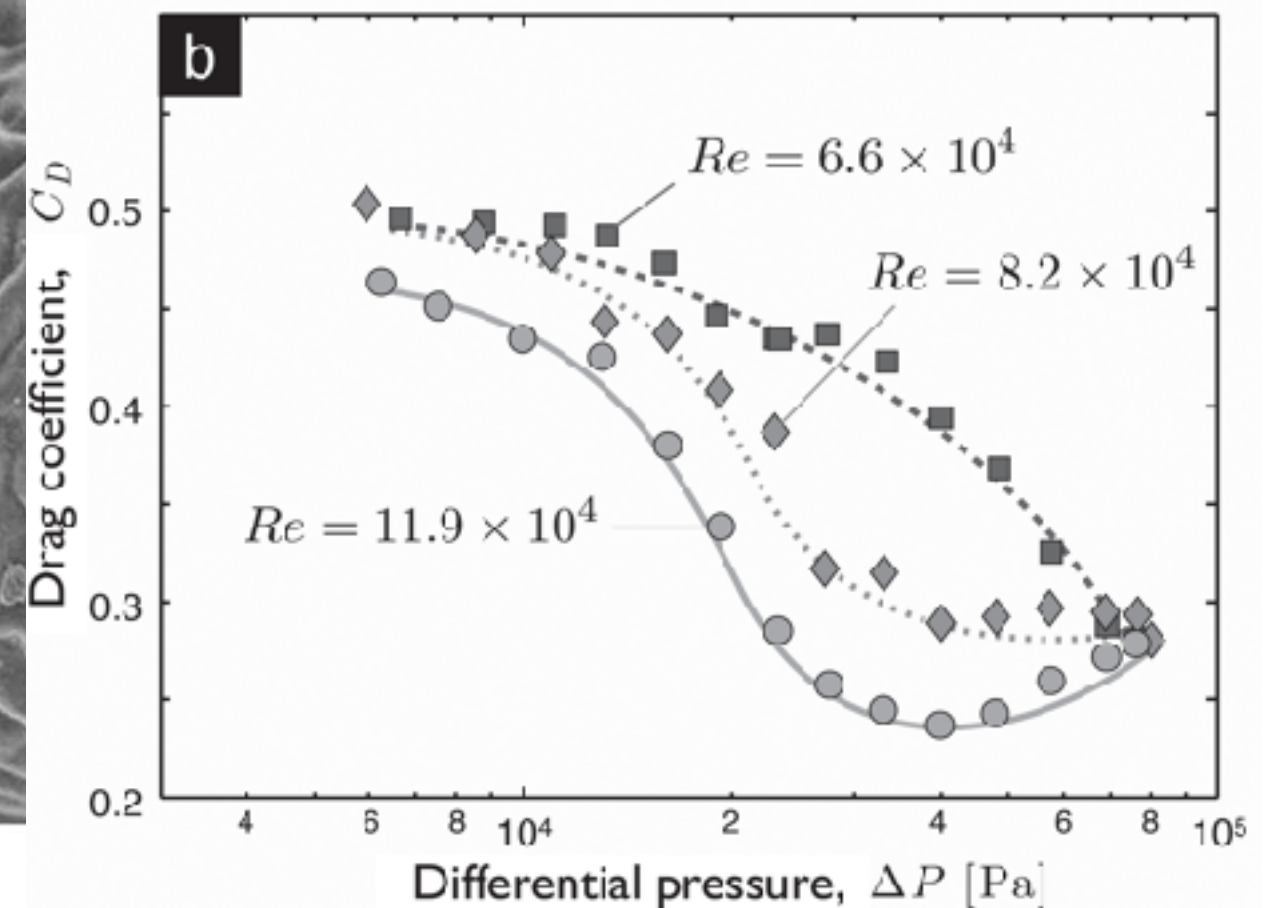
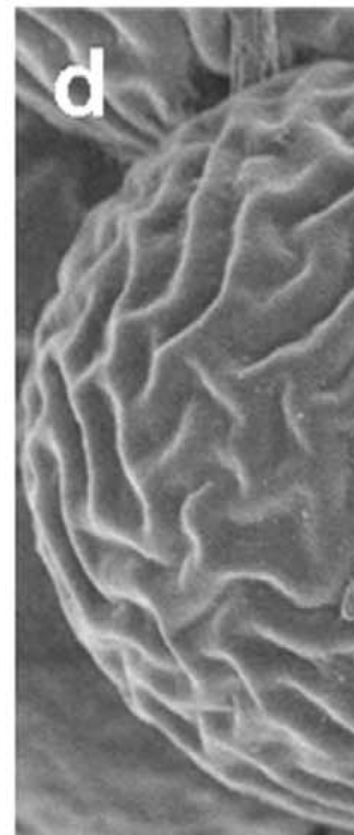
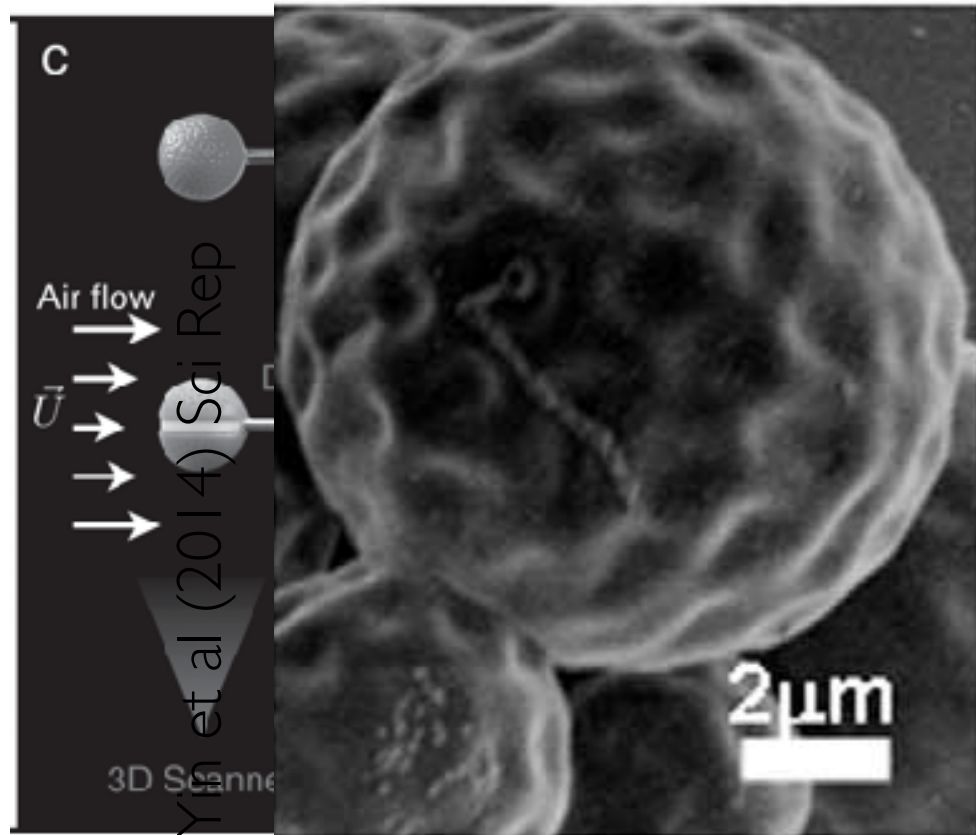
Curvature / stress-induced wrinkling transitions



by (2013) Soft Matter



$R = 805 \mu\text{m}$



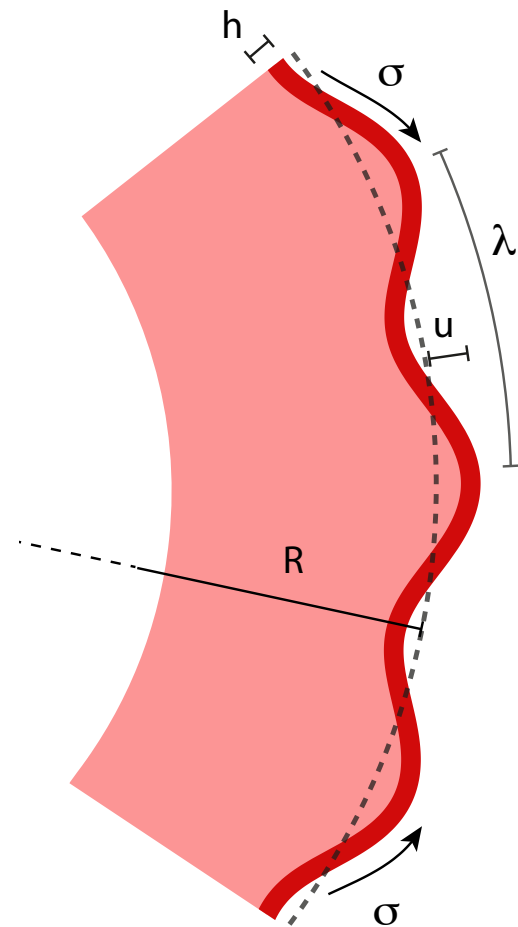
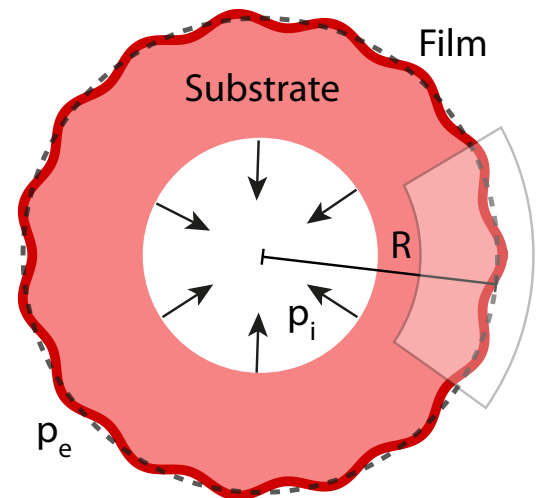
Generalized Swift-Hohenberg theory



Norbert Stoop

$$\partial_t u = \gamma_0 \Delta u - \gamma_2 \Delta^2 u - au - bu^2 - cu^3$$

$$+ \Gamma_1 [(\nabla u)^2 + 2u\Delta u] + \Gamma_2 [u(\nabla u)^2 + u^2 \Delta u]$$



Small deformations of a **sphere**

$$(a_{\alpha\beta}) = \begin{pmatrix} (R \sin \theta_2)^2 & 0 \\ 0 & R^2 \end{pmatrix}$$

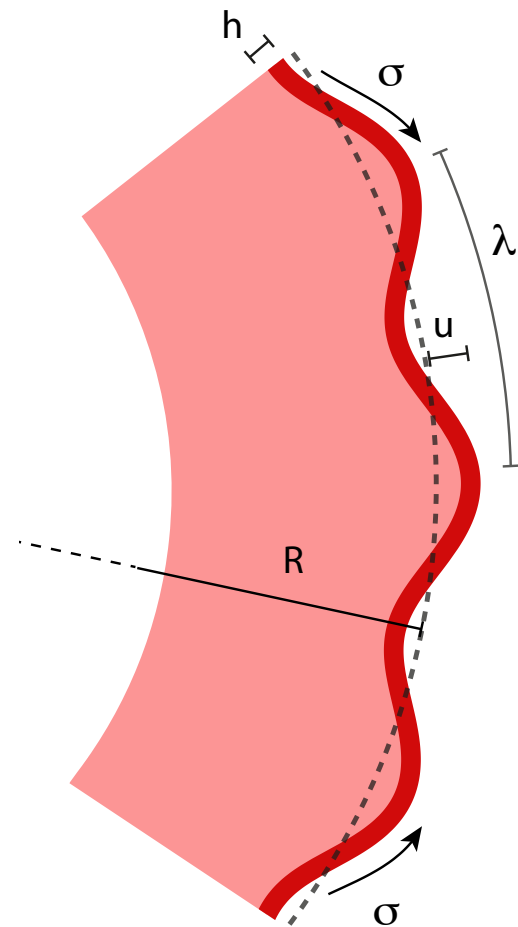
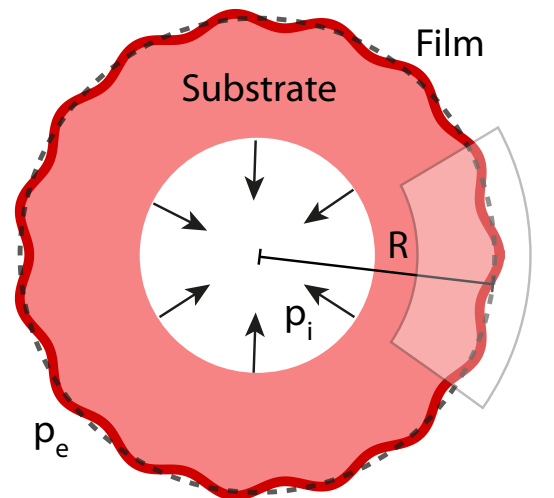
$$\Delta \psi = \nabla_\alpha \nabla^\alpha \psi = a^{\gamma\delta} \psi_{,\gamma\delta} - a^{\gamma\delta} \Gamma_{\gamma\delta}^\lambda \psi_{,\lambda}$$

Generalized Swift-Hohenberg theory



Norbert Stoop

$$\begin{aligned} \partial_t u = & \gamma_0 \Delta u - \gamma_2 \Delta^2 u - au - bu^2 - cu^3 \\ & + \Gamma_1 [(\nabla u)^2 + 2u\Delta u] + \Gamma_2 [u(\nabla u)^2 + u^2 \Delta u] \end{aligned}$$

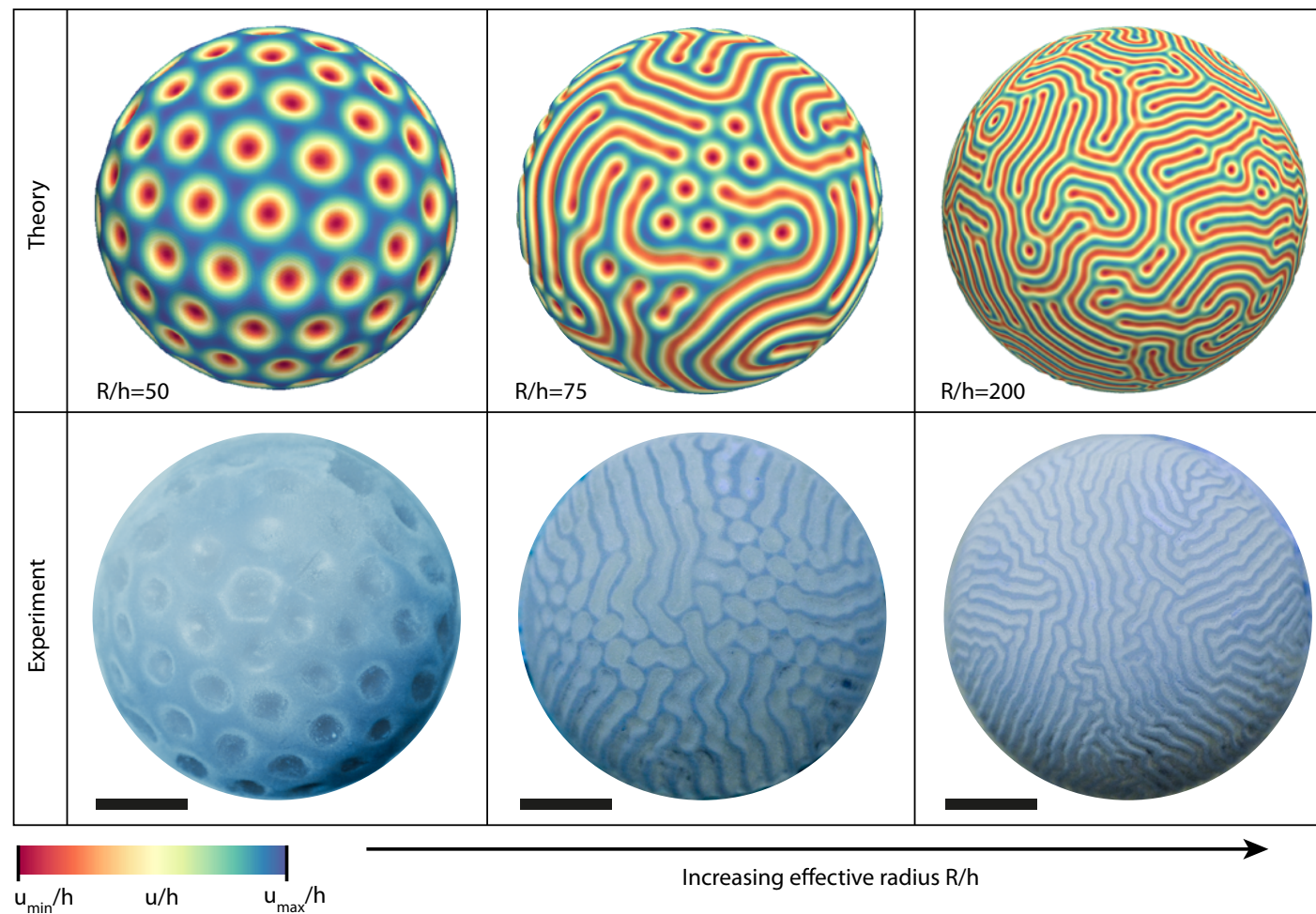


$$\begin{aligned} \gamma_0 &= \frac{\kappa^2}{3} - \frac{1}{6} \sqrt{\eta^{4/3} + 24(1+\nu)\kappa^2 + 16\kappa^4} \\ a &= \frac{\eta^{4/3}}{12} + \frac{6(1+\nu) - \eta^{2/3}}{3} \kappa^2 + \frac{\kappa^4}{3} + \tilde{a}_2 \Sigma_e \\ b &= 3(1+\nu)\kappa^3 \\ c &= \frac{2(1+\nu)\eta^{2/3}}{3} c_1 + (1+\nu)\kappa^4 \\ \Gamma_1 &= \frac{1+\nu}{2} \kappa \\ \Gamma_2 &= \frac{1+\nu}{2} \kappa^2 \\ \tilde{a}_2 &= -\frac{\eta^{4/3}(c + 3|\gamma_0|\Gamma_2)}{48\gamma_0^2} \end{aligned}$$

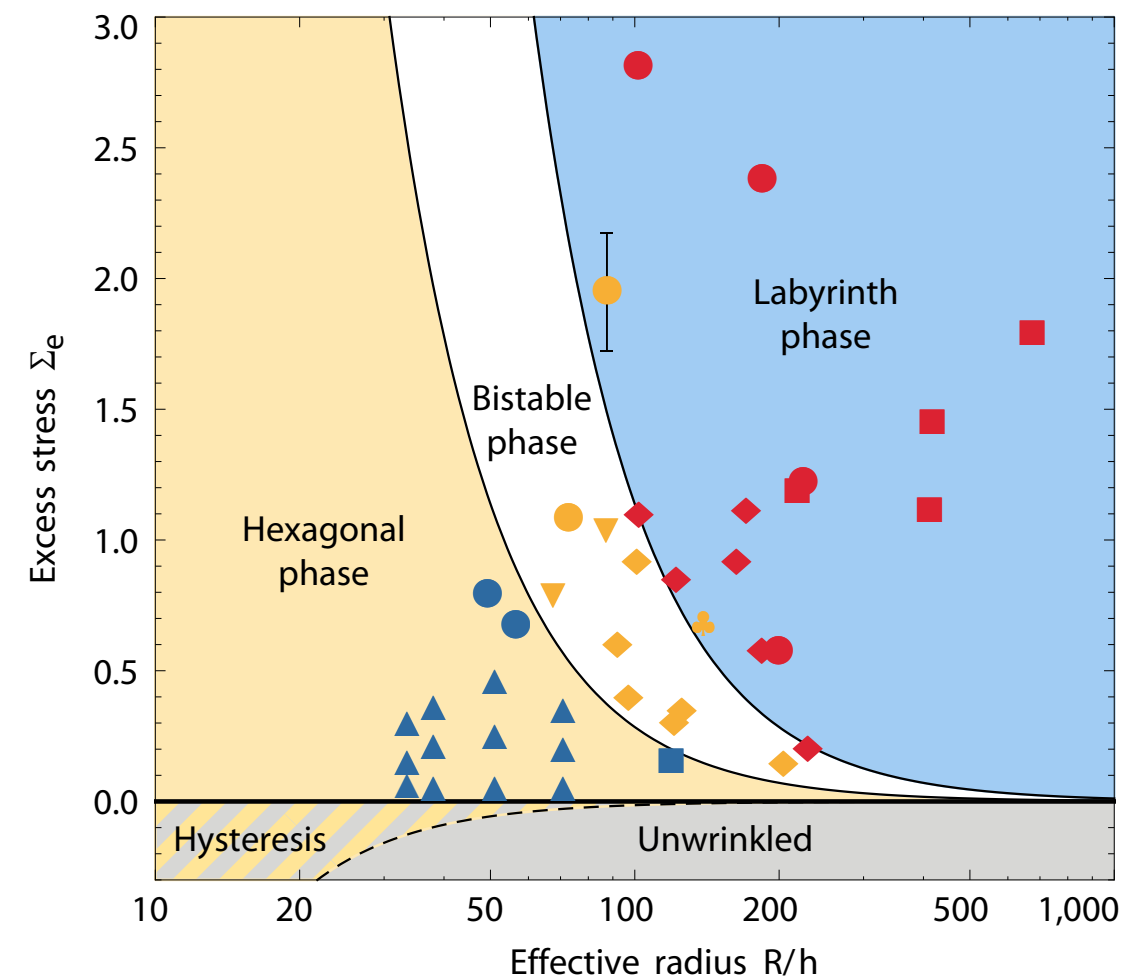
TABLE I: List of parameters for Eq. (1) in units $h = 1$, with $\eta = 3E_s/E_f$, $\gamma_2 = 1/12$, $\Sigma_e = (\sigma/\sigma_c) - 1$ and $\kappa = h/R$. The only remaining fitting parameter of the model is c_1 .

Theory correctly predicts

morphology



phase transitions



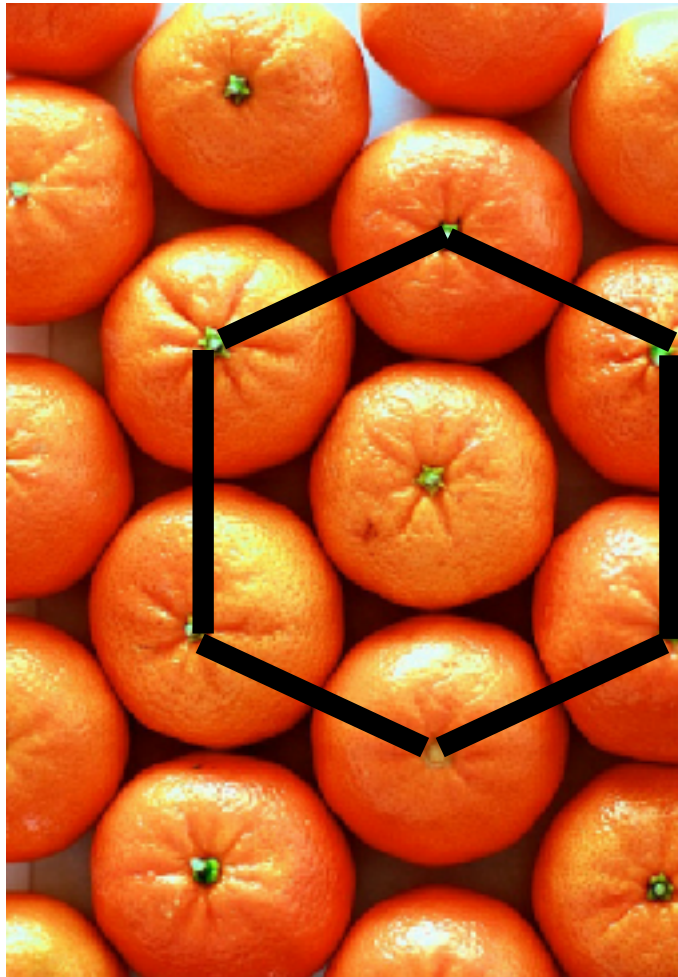
Arbitrarily curved surfaces

$$\begin{aligned} \mu \partial_t u = & \gamma_0 \Delta u - \gamma_2 \Delta^2 u - au - bu^2 - cu^3 + \\ & \frac{h}{2} \left\{ (\nu - 1) \left[b^{\alpha\beta} \nabla_\alpha u \nabla_\beta u + 2u \nabla_\beta (b^{\alpha\beta} \nabla_\alpha u) \right] + \right. \\ & \left. 2\nu \left[\mathcal{H}(\nabla u)^2 - 2\nabla \cdot (\mathcal{H}u \nabla u) \right] \right\} + \\ & \frac{h}{2} \left[(1 - \nu) u \nabla_\beta (uc^{\alpha\beta} \nabla_\alpha u) - \nu \mathcal{R}u (\nabla u)^2 + \right. \\ & \left. \nu \nabla \cdot (\mathcal{R}u^2 \nabla u) \right] \end{aligned}$$

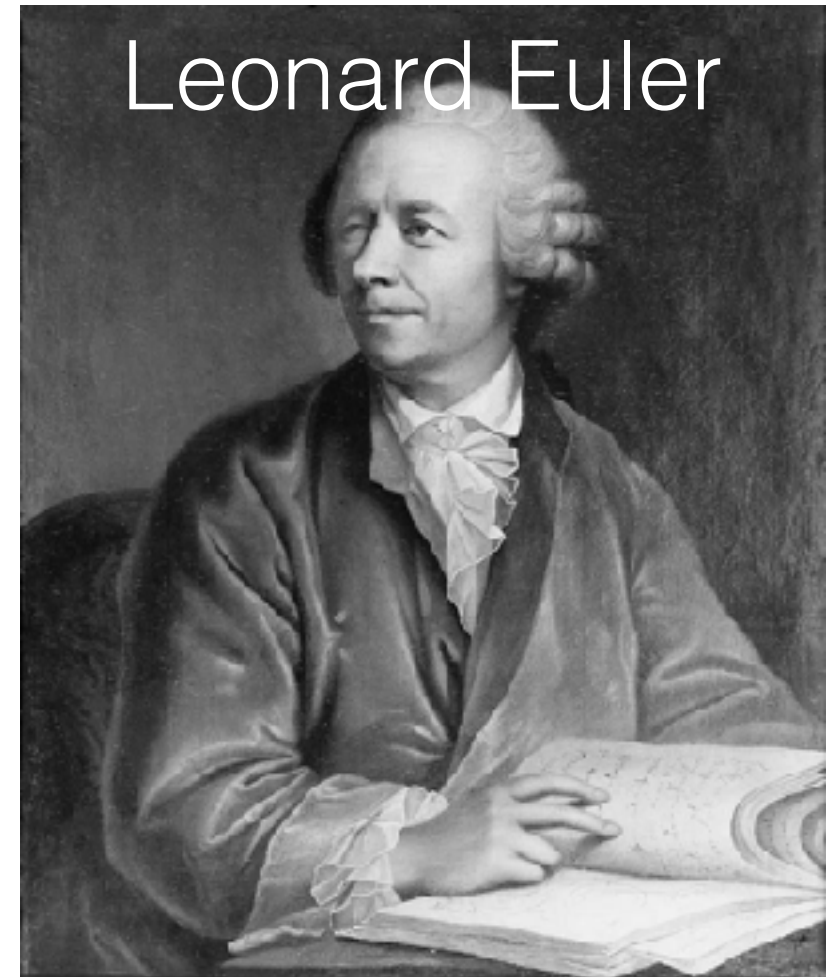




Topological defects



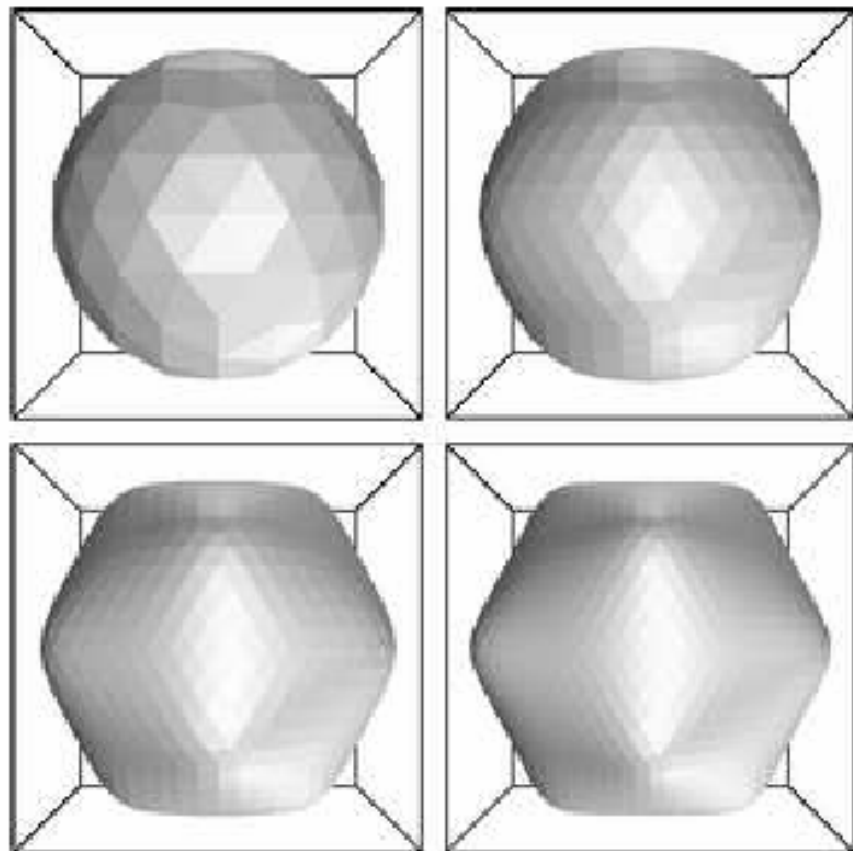
20 hexagons
12 pentagons



$$V - E + F = 2$$

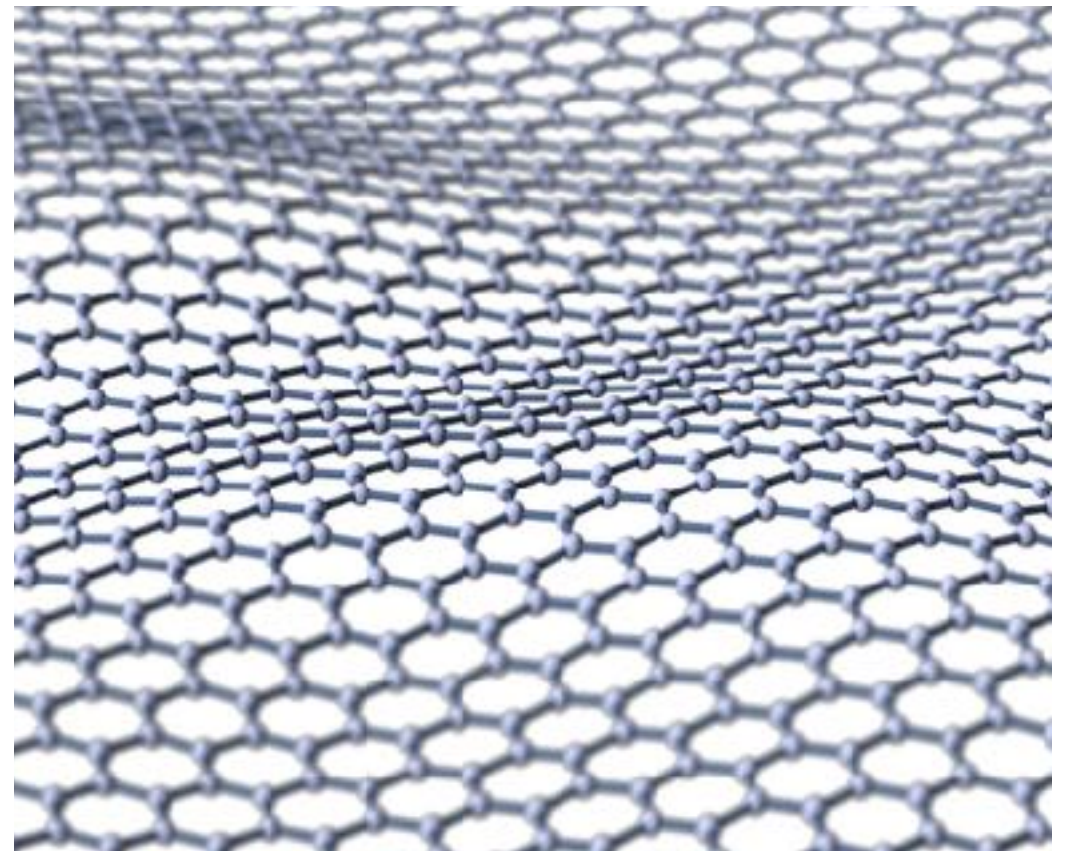
Why interesting ?

topological defects
nucleate size-induced shape transition
in viral capsids



J Lidmar, L Mirny and DR Nelson (2003) PRE

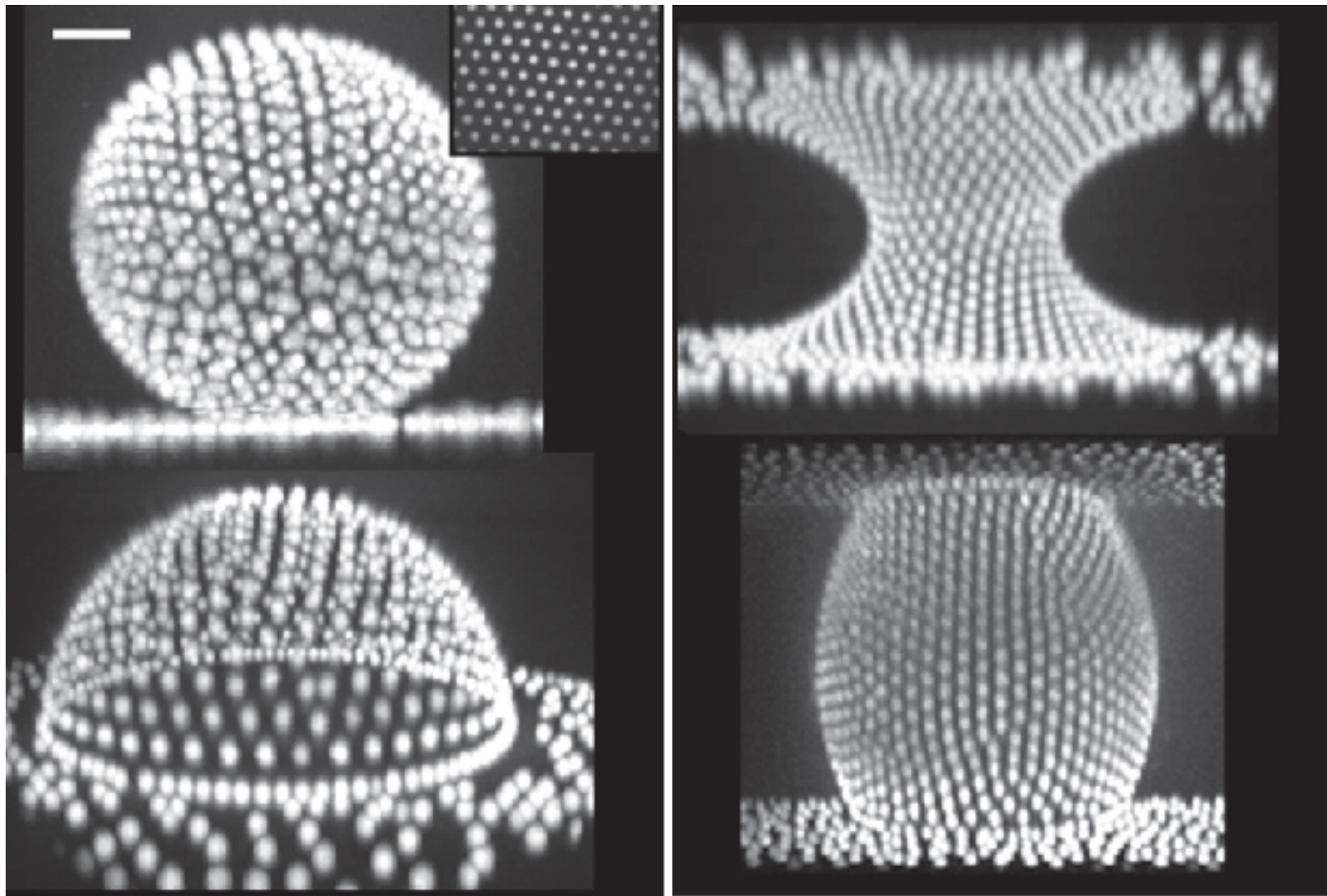
bending of graphene
introduces defects and
changes electronic properties



A Cortijo and MAH Vozmediano (2007) EPL

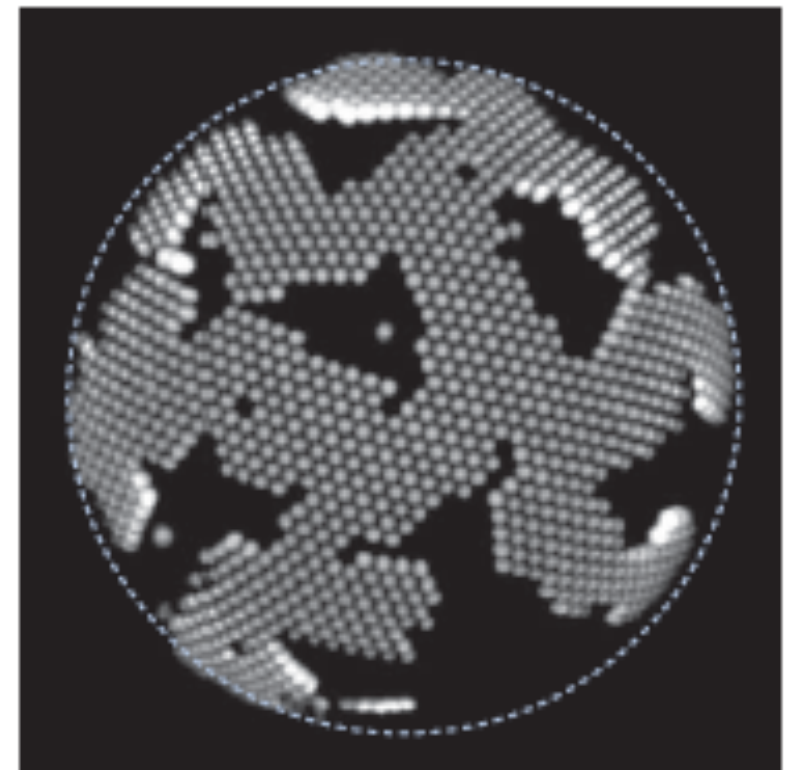
Surface crystallography

Statics



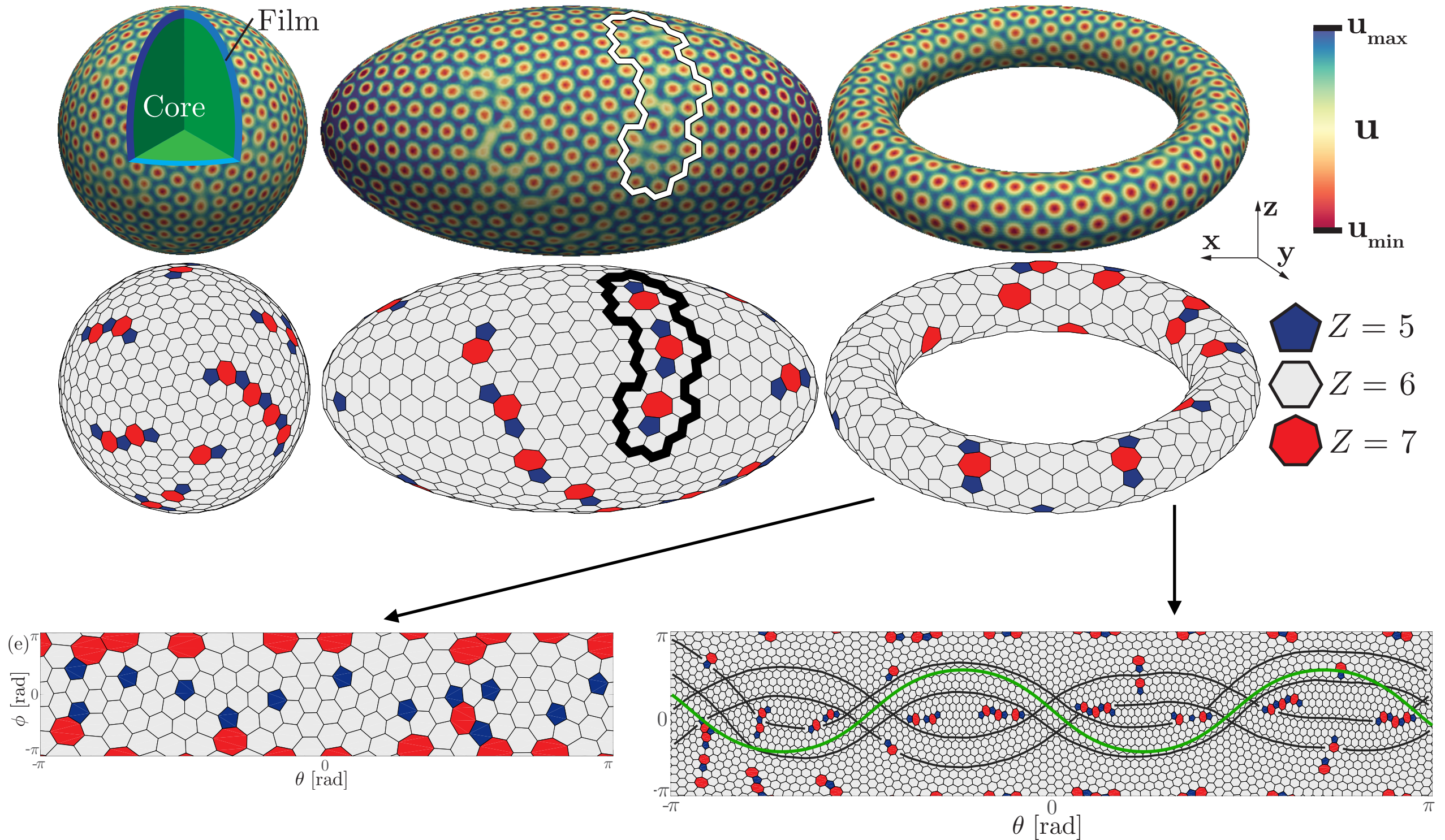
Irvine et al (2010) Nature

Nucleation



Meng, Paulose, Nelson & Manoharan
(2014) Science

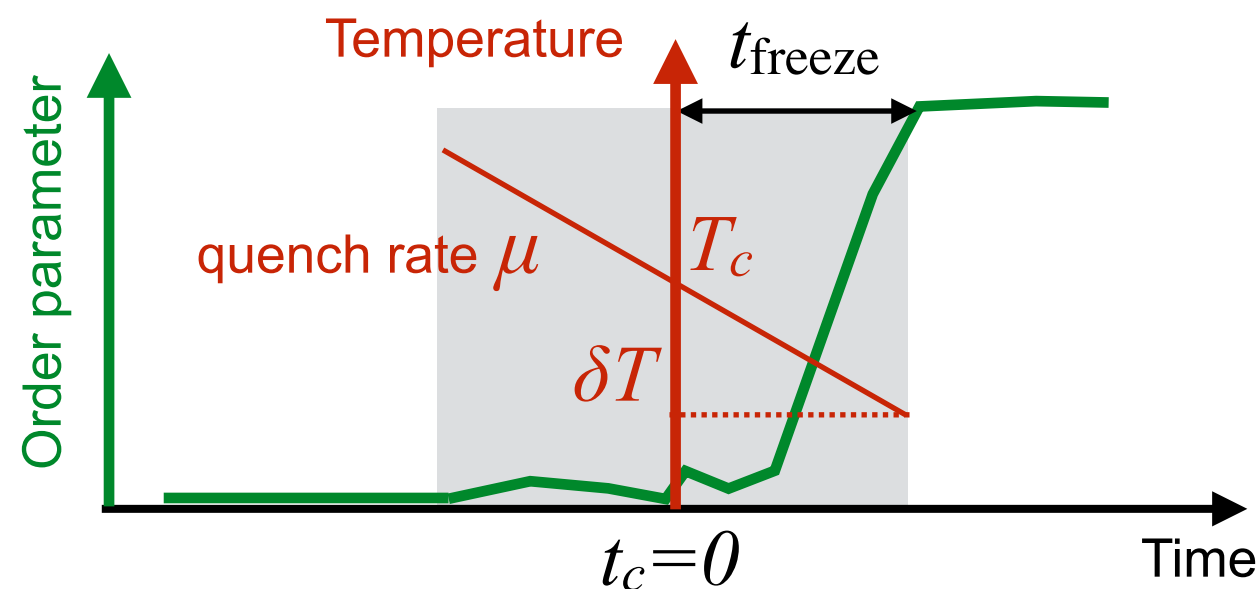
Statics: surface crystallography



Dynamics: Kibble-Zurek mechanism (KZM)

Kibble & Zurek (1970s): System driven through a 2nd order phase transition

- exhibits critical slowing-down
- dynamics cannot follow changes of external system parameters
- density of topological defects after quench reveals information about the quench dynamics
- observed topological structures in the universe provide a window into early evolution

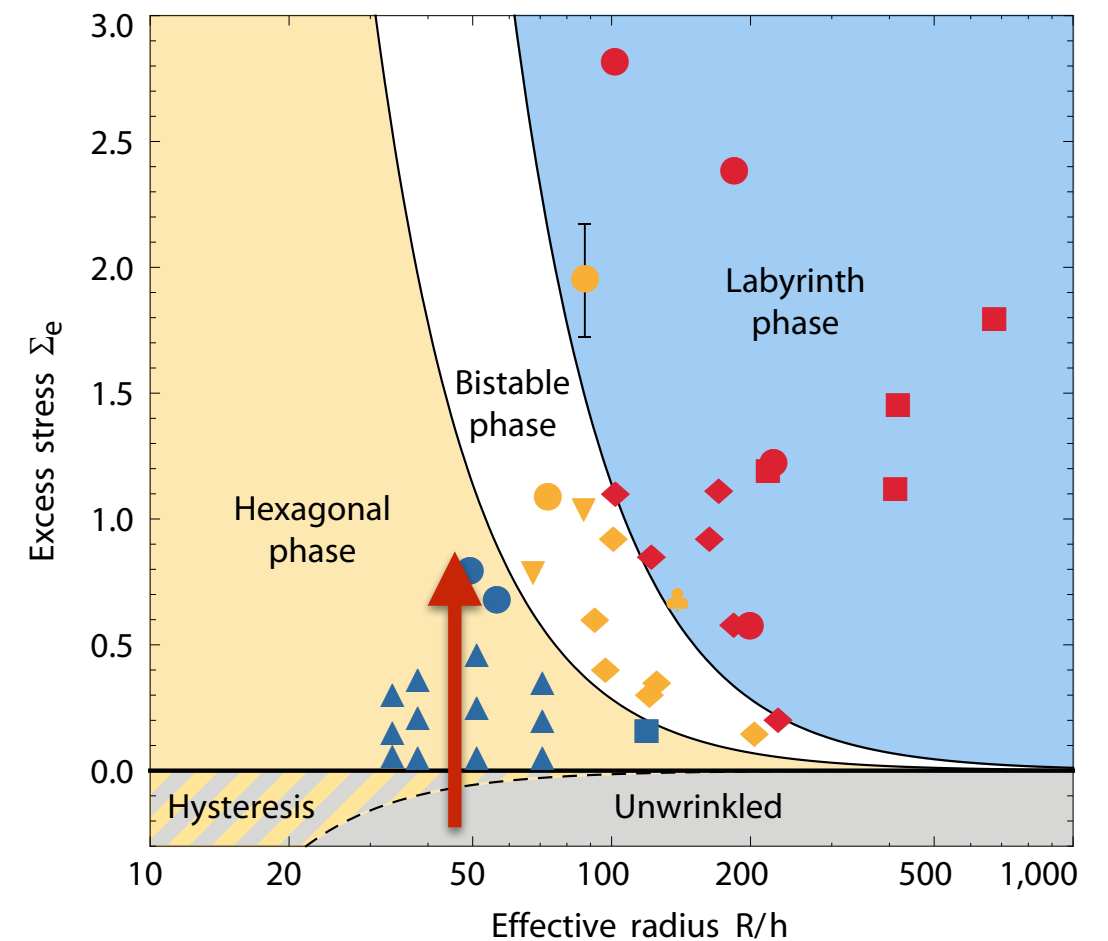


KZ predictions:

$$T(t) = T_c - \mu t \Rightarrow t_{\text{freeze}} \sim \mu^{-1/2}$$

$$\delta T(t_{\text{freeze}}) \sim \mu^{1/2}$$

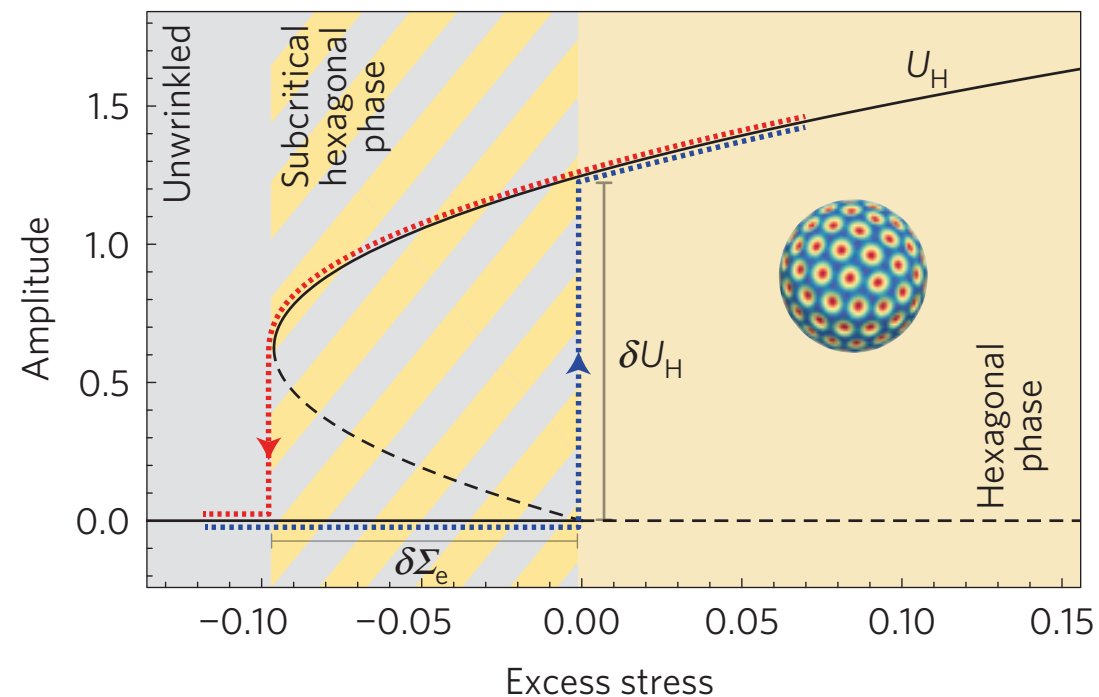
$$\rho_{\text{defect}} \sim \mu^{1/2}$$



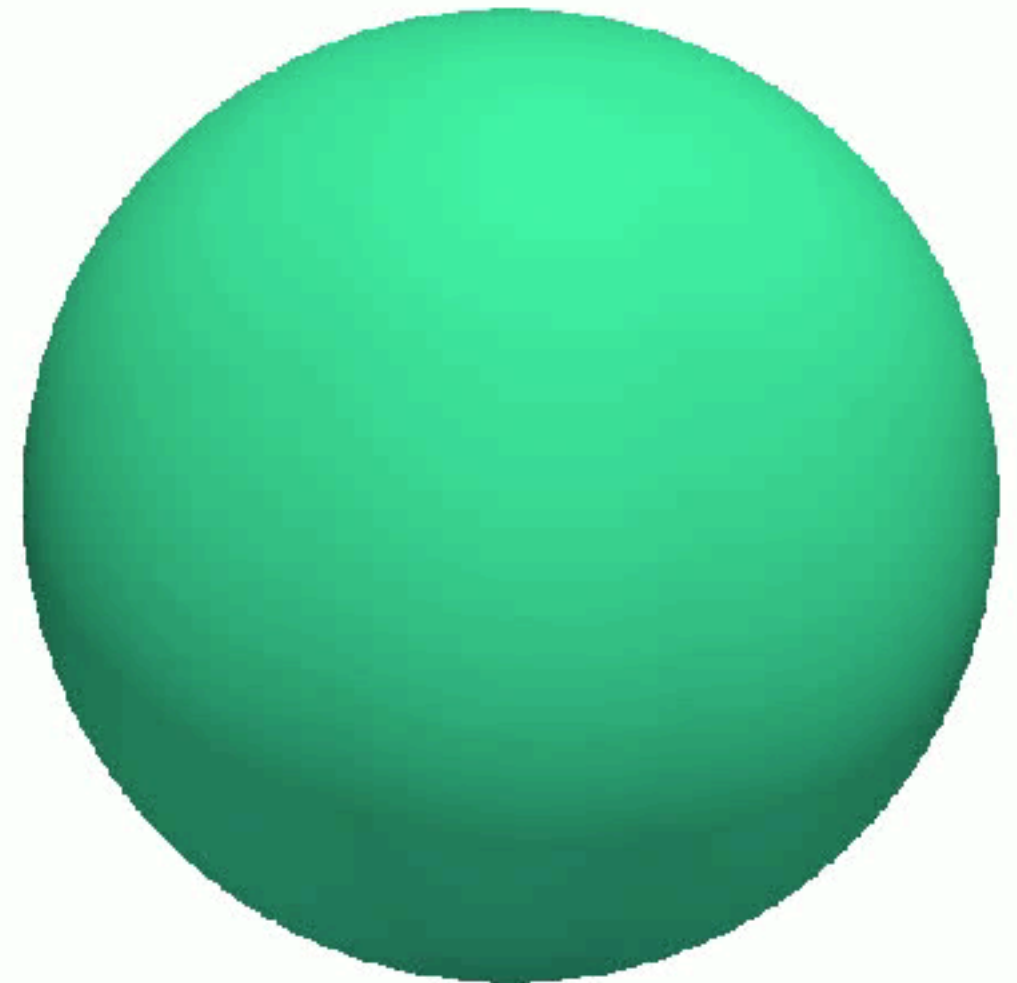
Elastic surface crystals as testbed for KZM in curved geometries?

Dynamics of phase transition

Adiabatic/equilibrium

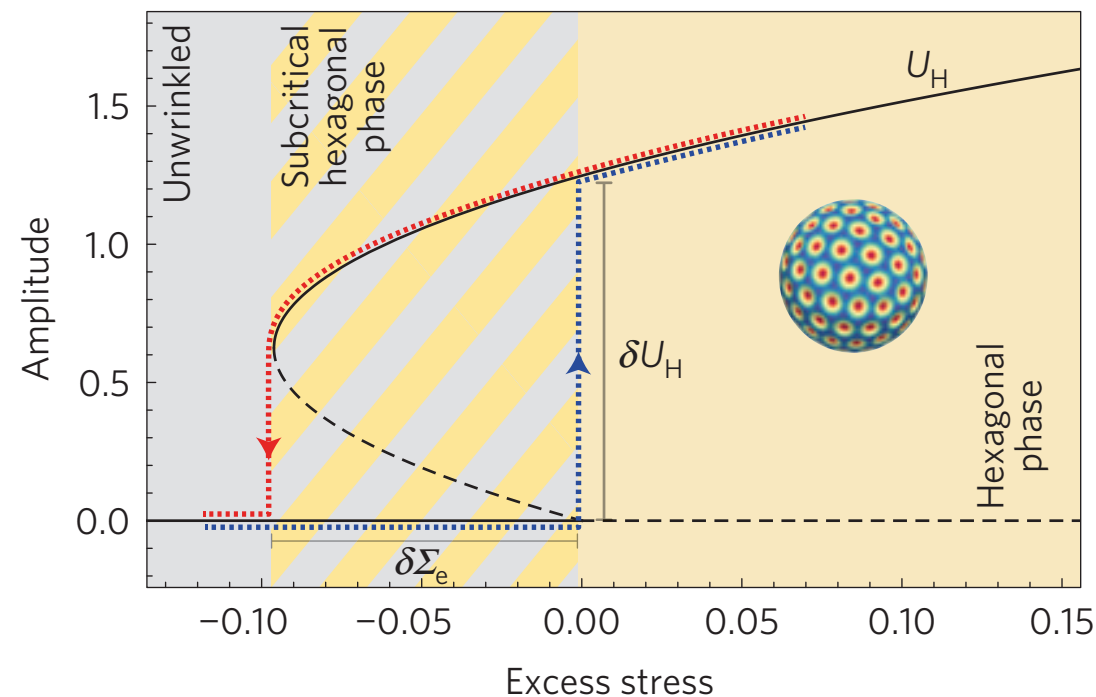


bifurcation from flat state $u=0$
to hexagonal pattern at $\Sigma_e = 0$



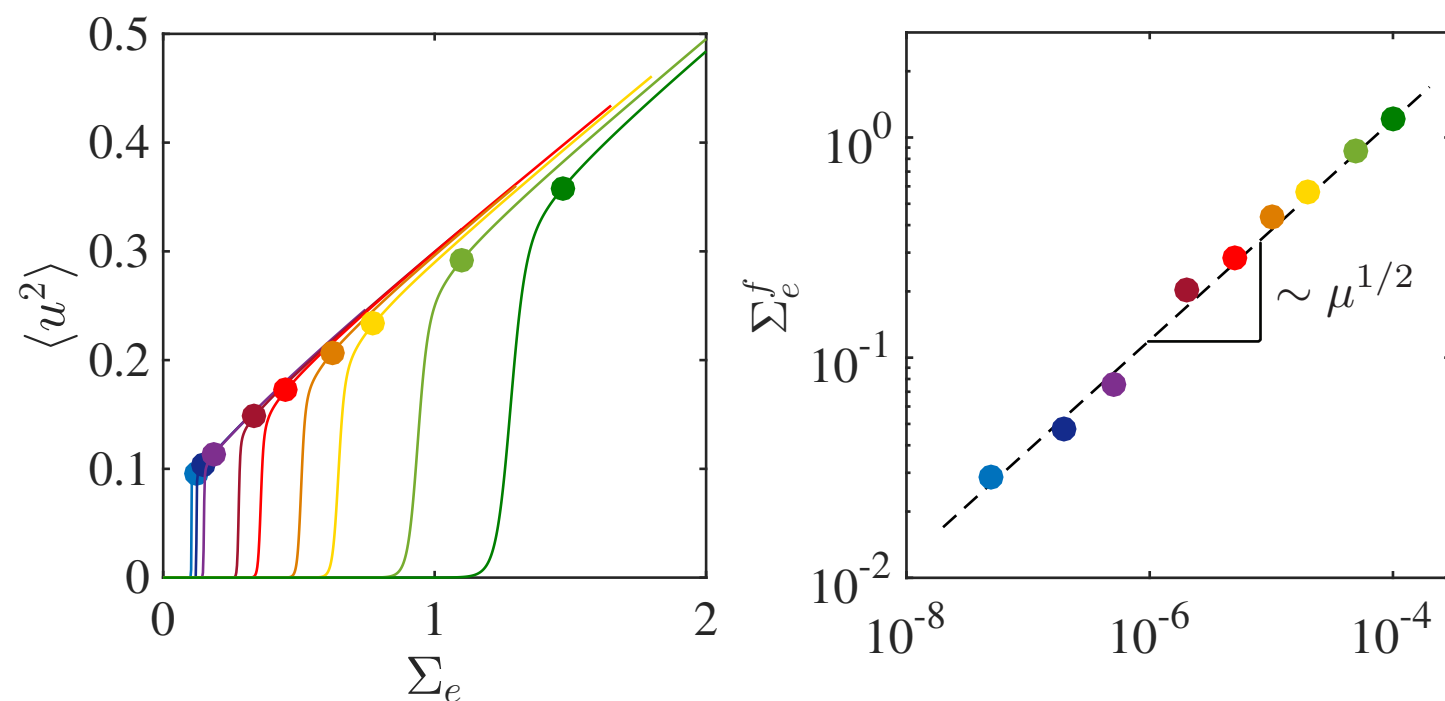
Freeze-out time follows KZ scaling

Adiabatic/equilibrium



bifurcation from flat state $u=0$
to hexagonal pattern at $\Sigma_e=0$

Linear quench $\Sigma_e(t) = \mu t$



$$\Sigma_e^f \sim \mu^{1/2}$$

Freeze-out time follows KZ scaling

$$u(t, \mathbf{x}) = U(t) \sum_{a=1}^3 (e^{i\mathbf{k}_a \cdot \mathbf{x}} + e^{-i\mathbf{k}_a \cdot \mathbf{x}})$$

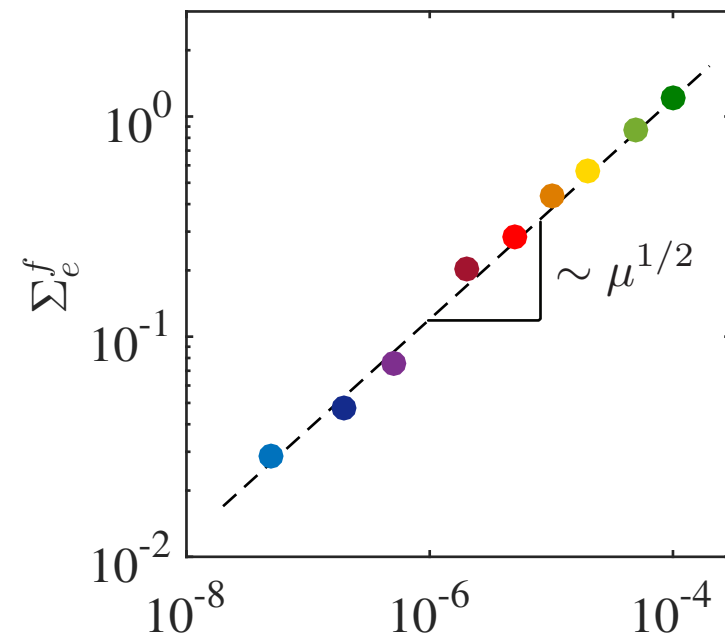
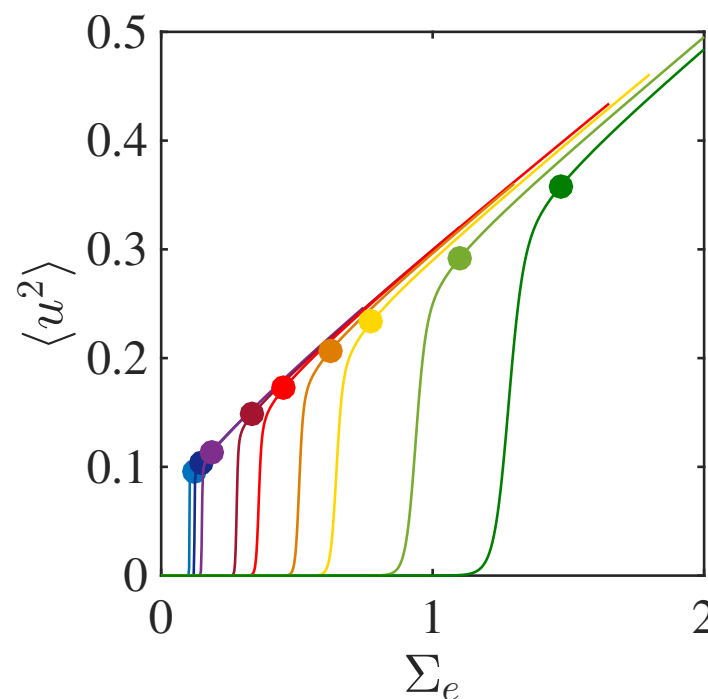
$$\frac{dU}{dt} = \frac{c\Sigma_e}{12\gamma_0^4}U - \frac{4}{3R|\gamma_0|^3}U^2 - \frac{15c}{9\gamma_0^4}U^3$$

$$\frac{d}{dt}U \approx \frac{c\mu t}{12\gamma_0^4}U$$

$$\frac{d}{dt'}U \approx \frac{ct'}{12\gamma_0^4}U$$

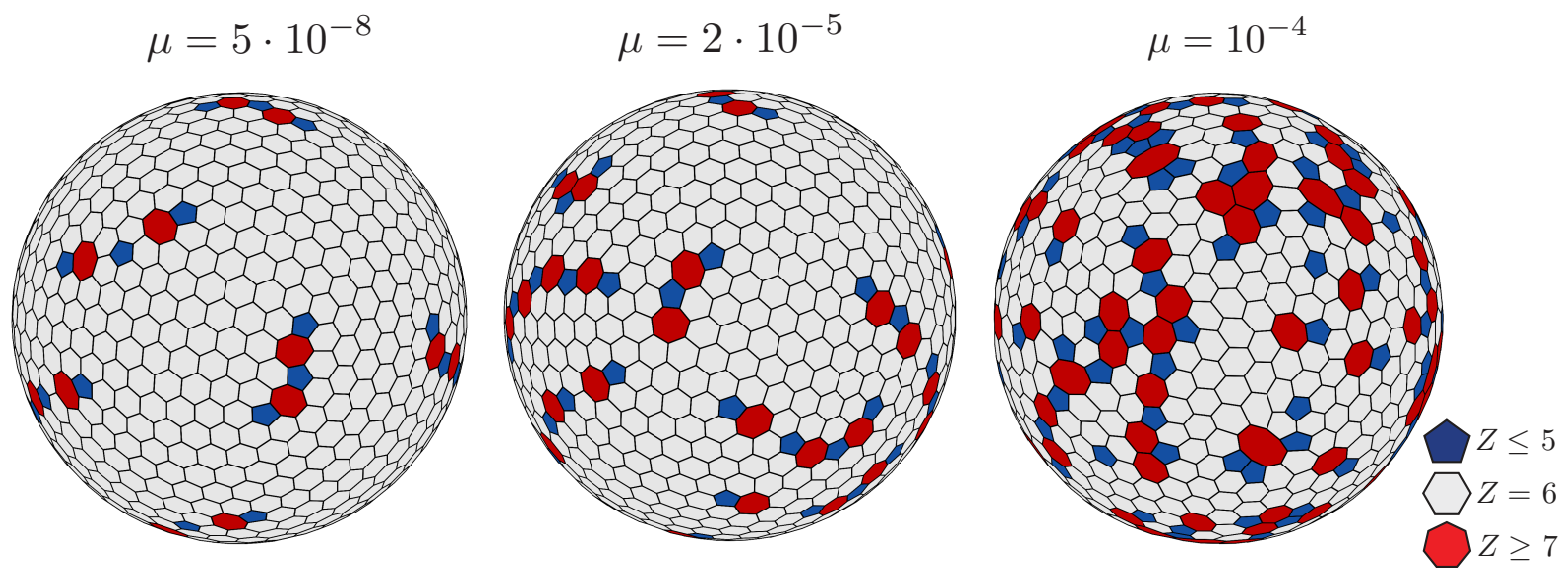
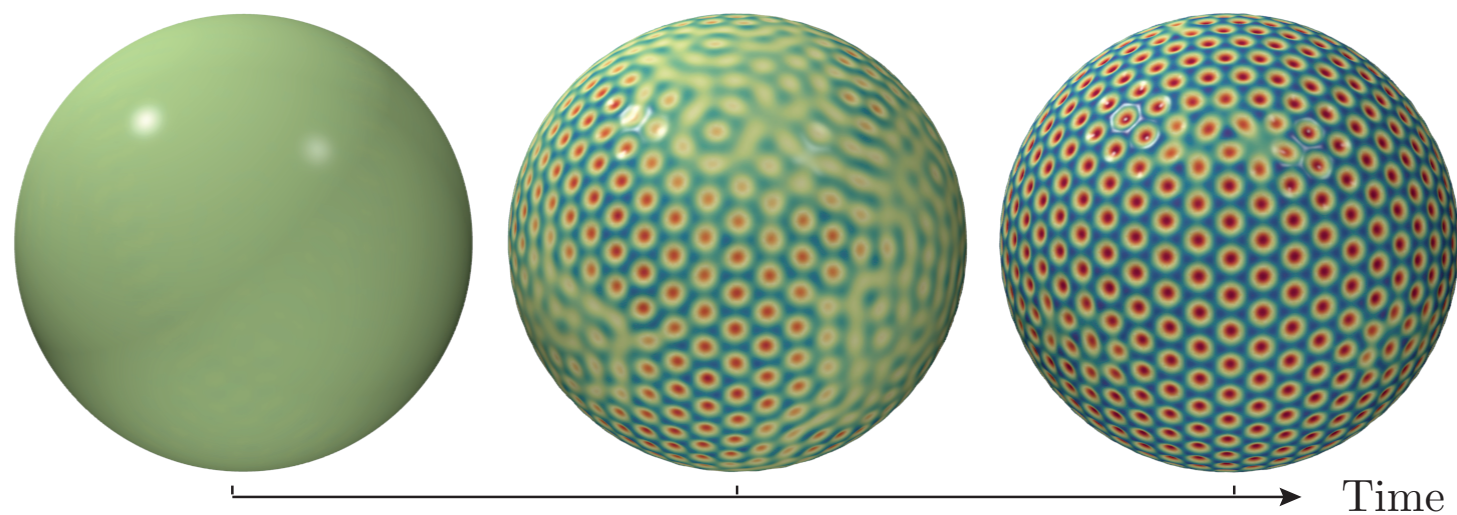
$$t' = \mu^{1/2}t$$

Linear quench $\Sigma_e(t) = \mu t$

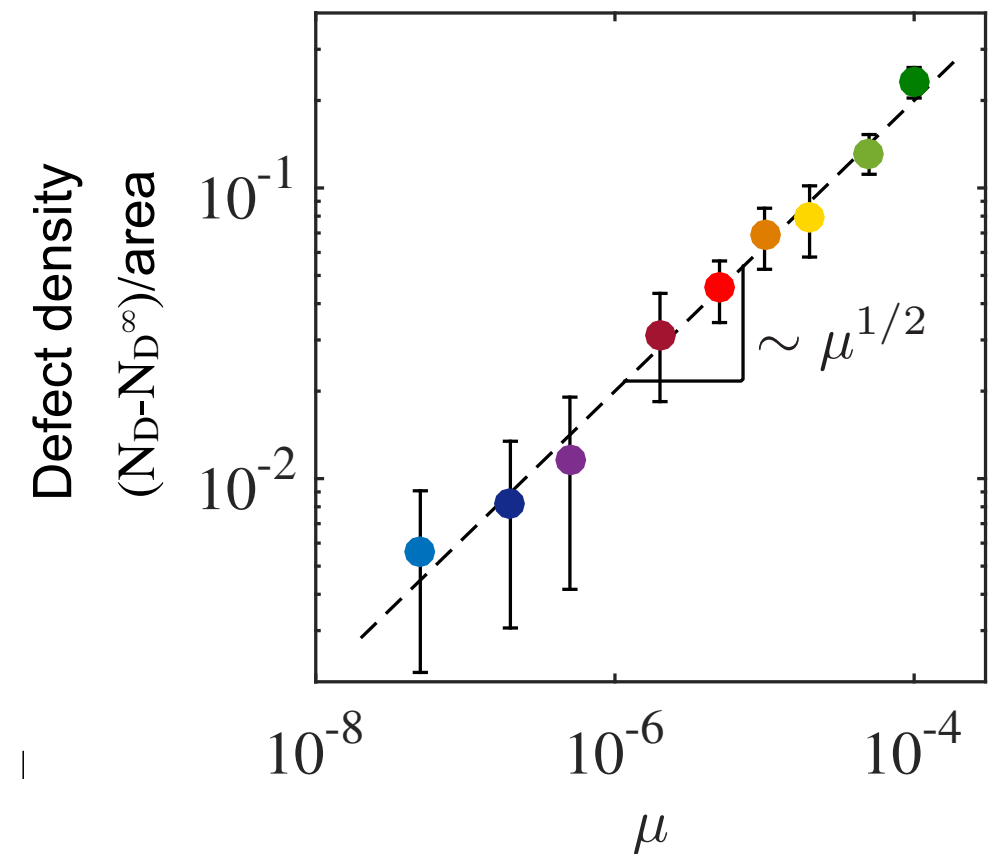


$$\Sigma_e^f \sim \mu^{1/2}$$

Defect density follows KZ predictions !?

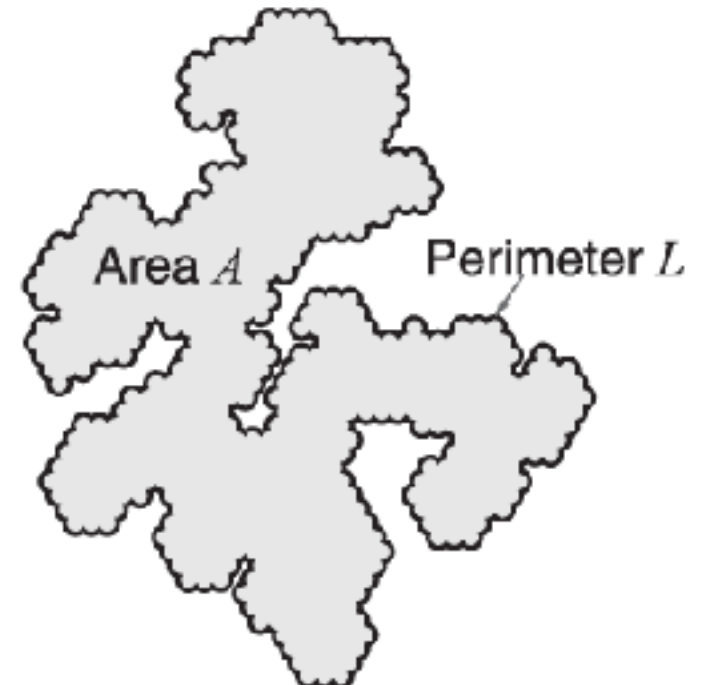
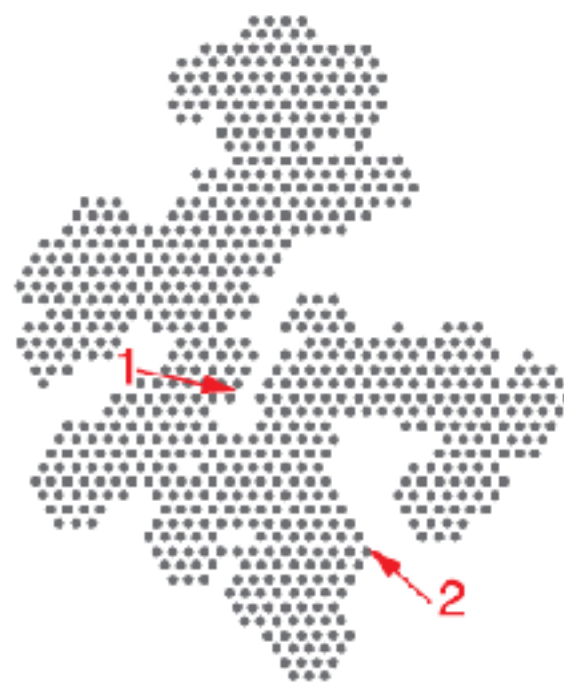
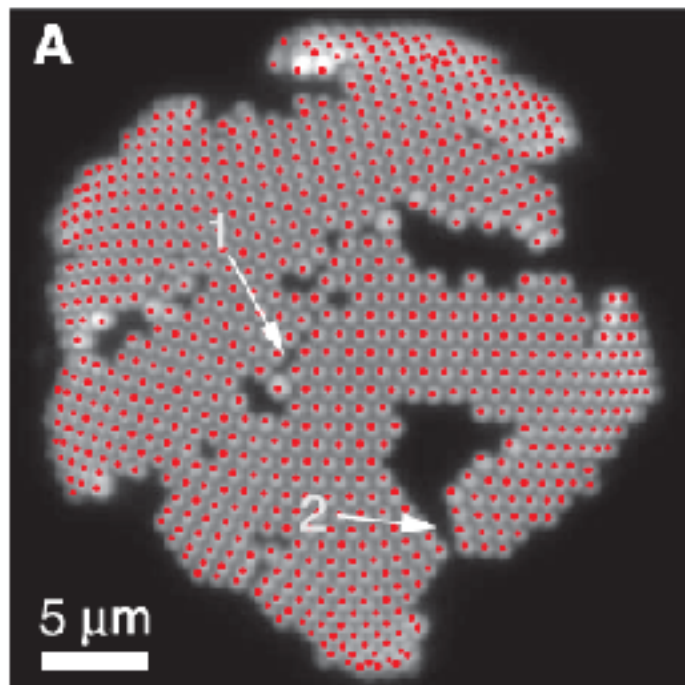
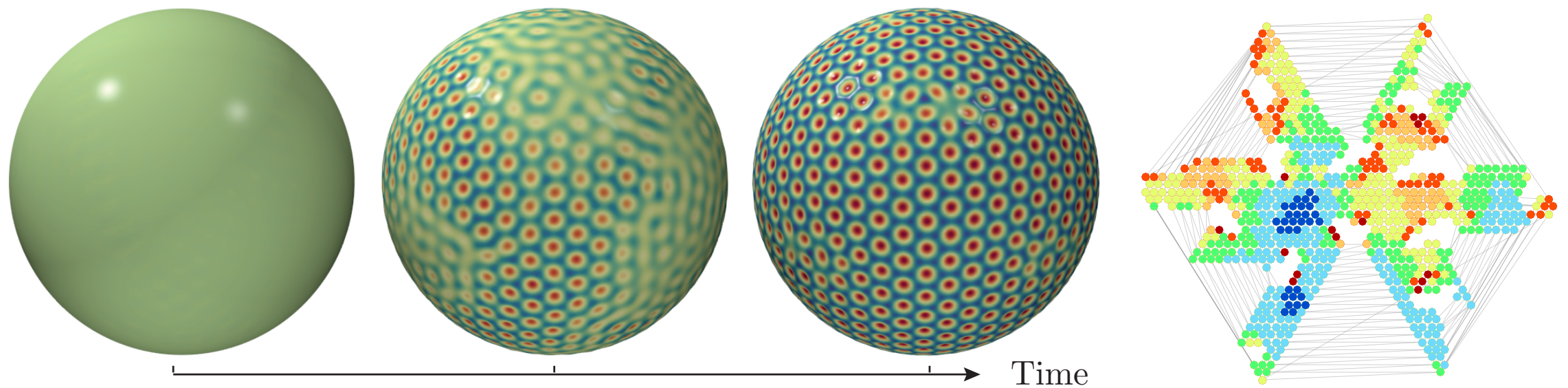


Voronoi tessellation at freeze-out Σ_e^f

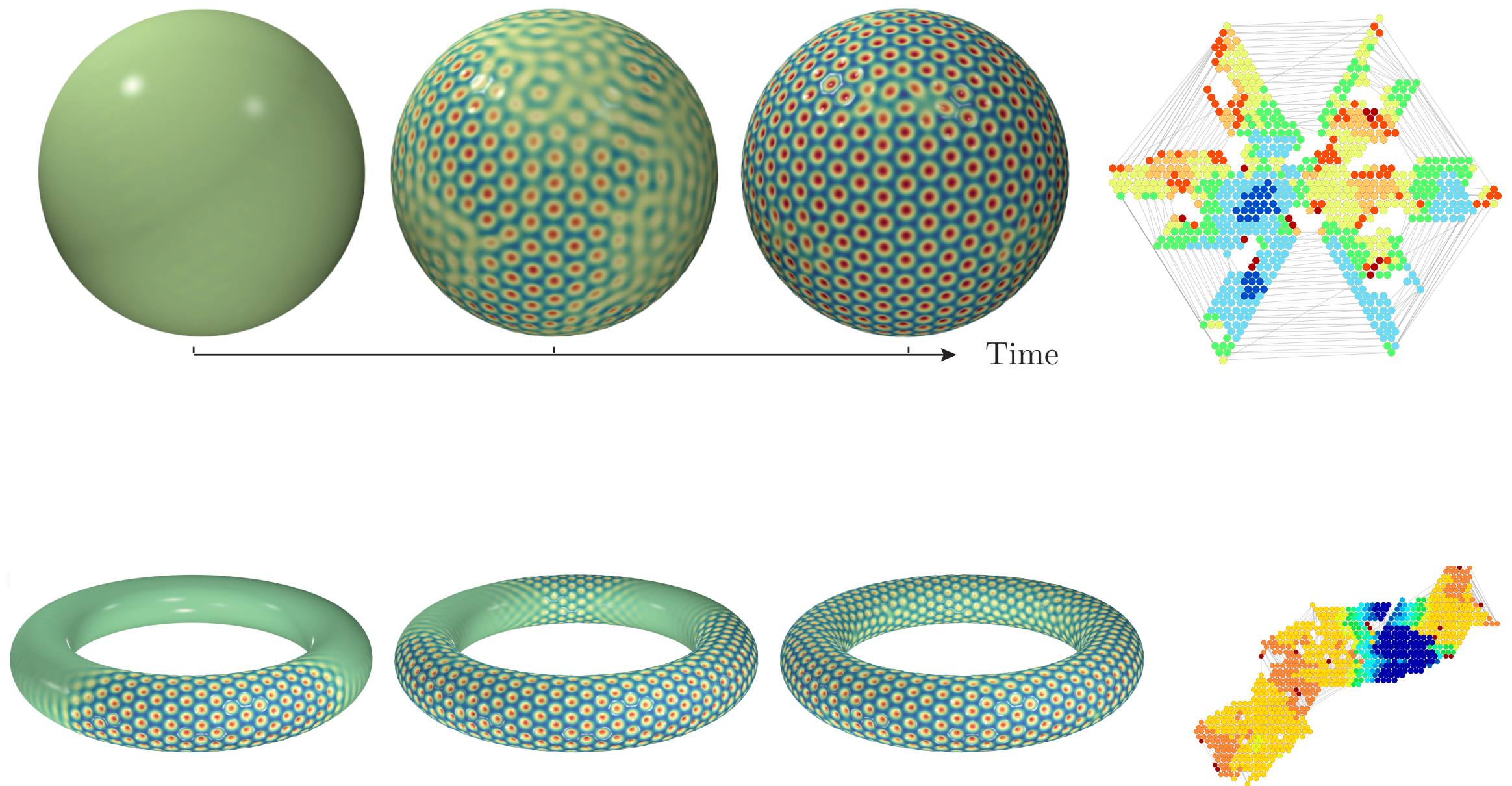


$$\rho_{\text{defect}} \sim \mu^{1/2}$$

Nucleation dynamics



Nucleation dynamics



explains 'geodesic wrapping'

